

Lagrangian coherent structures in fluid dynamics

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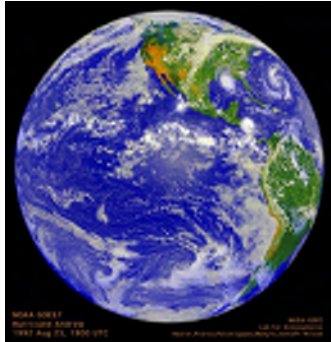
SES, October 2011



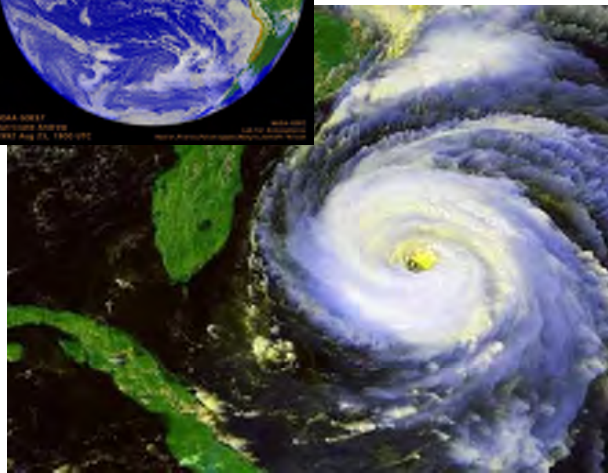
MultiSTEPS: MultiScale Transport in
Environmental & Physiological Systems,
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Lagrangian coherent structures – coherent structures moving with the fluid



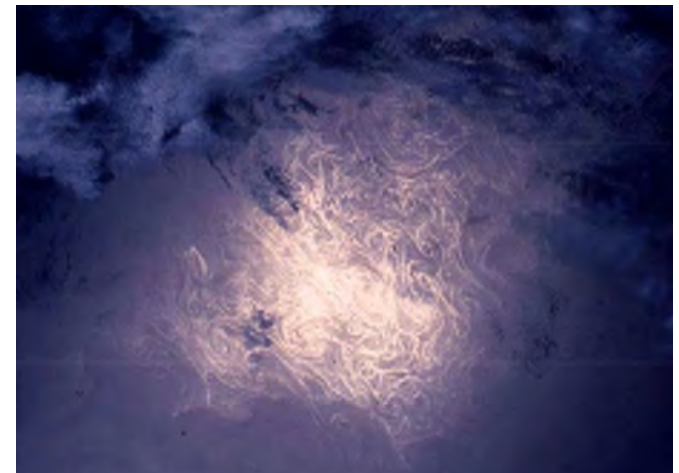
Some vortex examples from geophysical fluids



hurricanes



tornados



eddies

Why Lagrangian coherent structures?

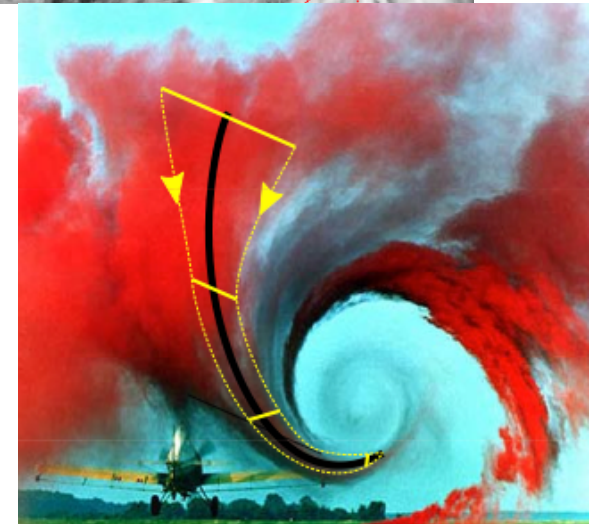
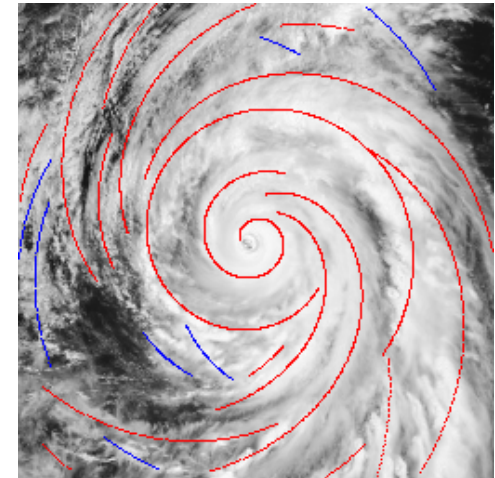
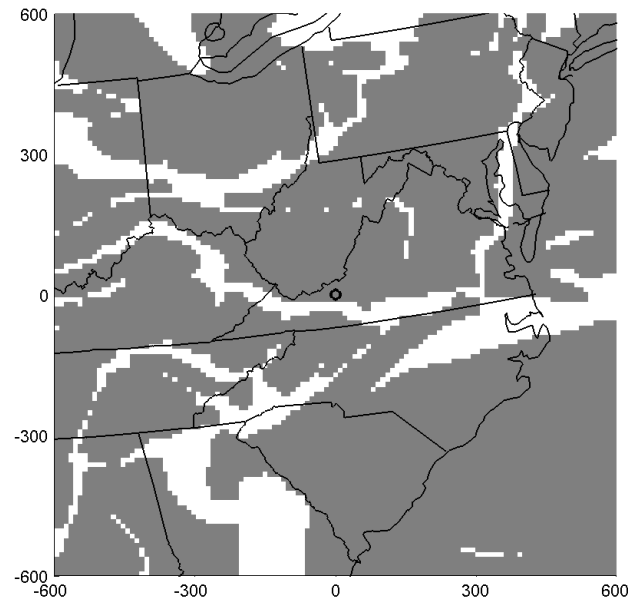
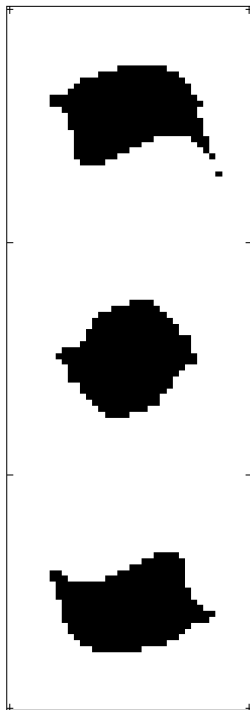
Natural flow visualization

Insight for design and control

Application to fluid and non-fluid systems

Lagrangian coherent structures

- Two strategies from *dynamical systems*:
- (1) Identify their *'skeleton'* or *boundaries*
 - (2) identify the structures *directly*



Lagrangian coherent structures

Two strategies from *dynamical systems*:

(1) Identify their *'skeleton'* or *boundaries*

(2) identify the structures *directly*

For (1):

For time-periodic or steady fluid velocity fields, identify periodic orbits + their stable & unstable manifolds.

But for *realistic* flows,

Time-independent and aperiodic, data-driven, finite-time, so need something else?

For (2):

Discretize the *flow map* over timescale of interest; determine coherent regions via eigenmodes, graph methods

Atmosphere: Antarctic polar vortex

Atmosphere: Antarctic polar vortex

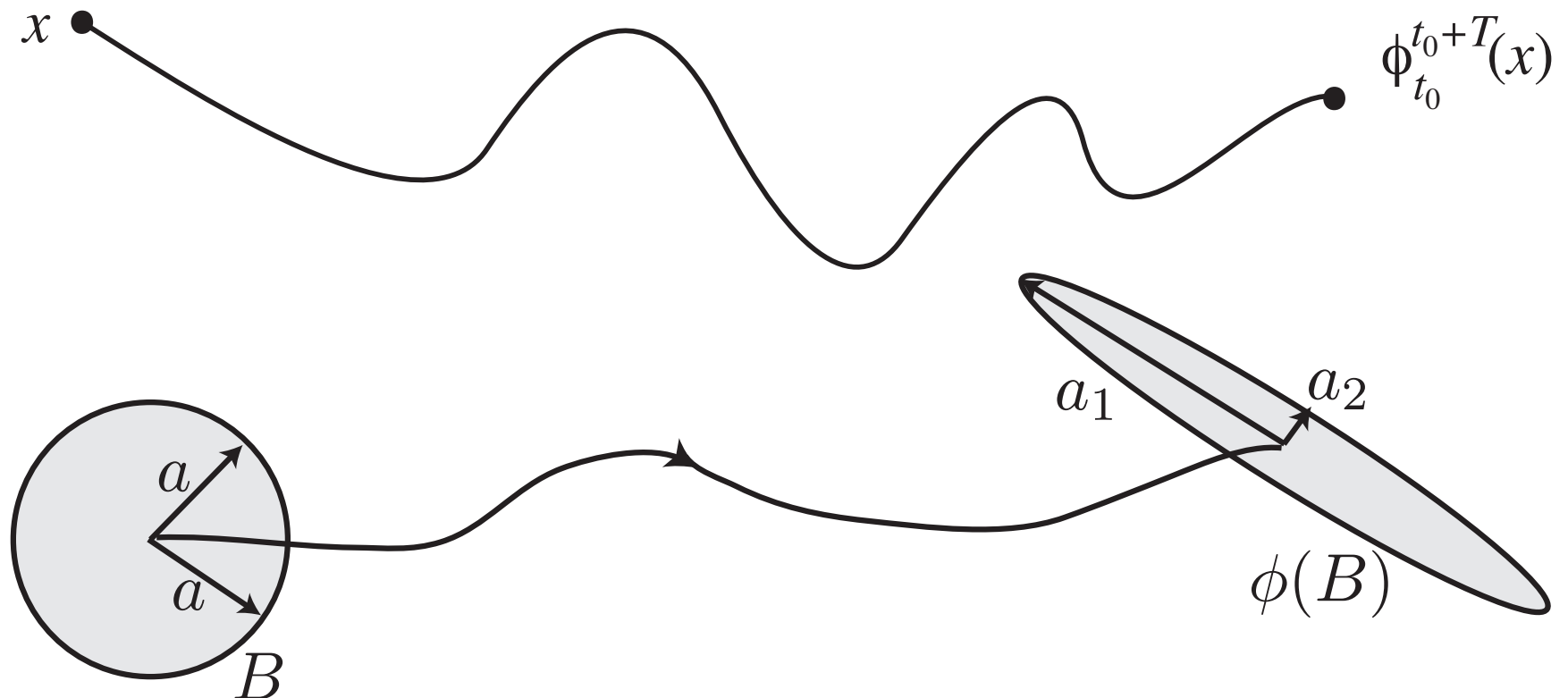
Atmosphere: continental U.S.

Periodic velocity field

- If \mathcal{M} = fluid domain, the flow map,

$$\phi_t^{t+T} : \mathcal{M} \longrightarrow \mathcal{M},$$

takes points $x \mapsto \phi_t^{t+T}(x)$ to their location after time T

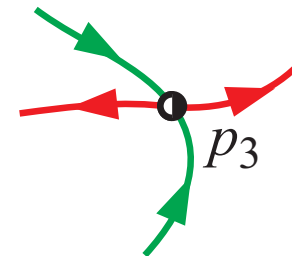
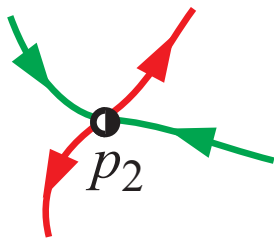
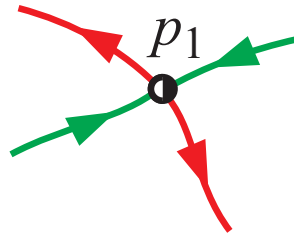


Periodic velocity field

- Suppose our velocity field is periodic with period T and we consider the flow map $f = \phi_t^{t+T}$ over one period
- To understand the transport of points under the map f , consider **invariant manifolds of saddle fixed points**
- Let $p_i, i = 1, \dots, N_p$, denote a collection of saddle-type hyperbolic fixed points for f .

Partition phase space into coherent regions

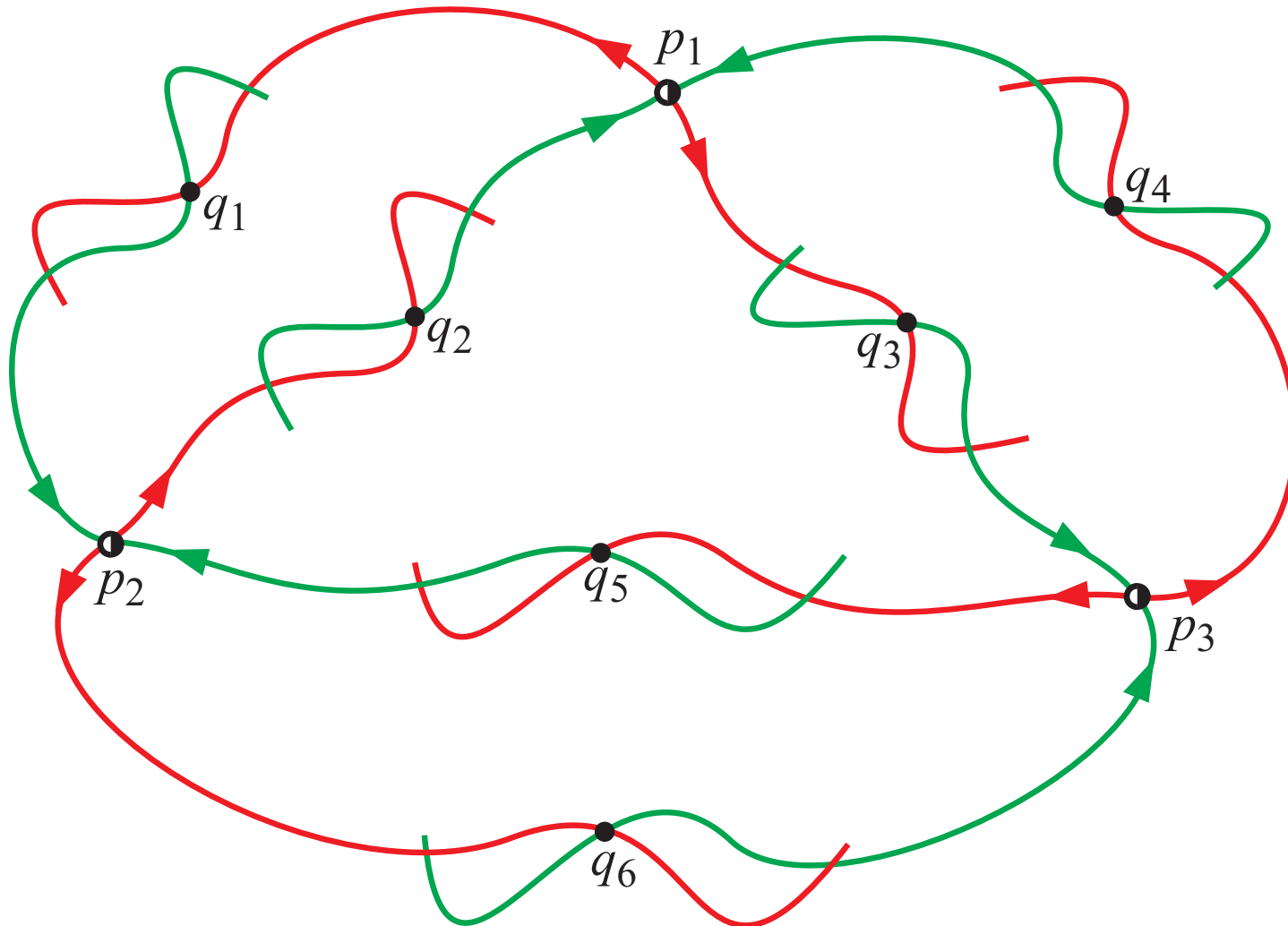
- Natural way to partition phase space
 - Pieces of $W^u(p_i)$ and $W^s(p_i)$ partition \mathcal{M} .



Unstable and stable manifolds in **red** and **green**, resp.

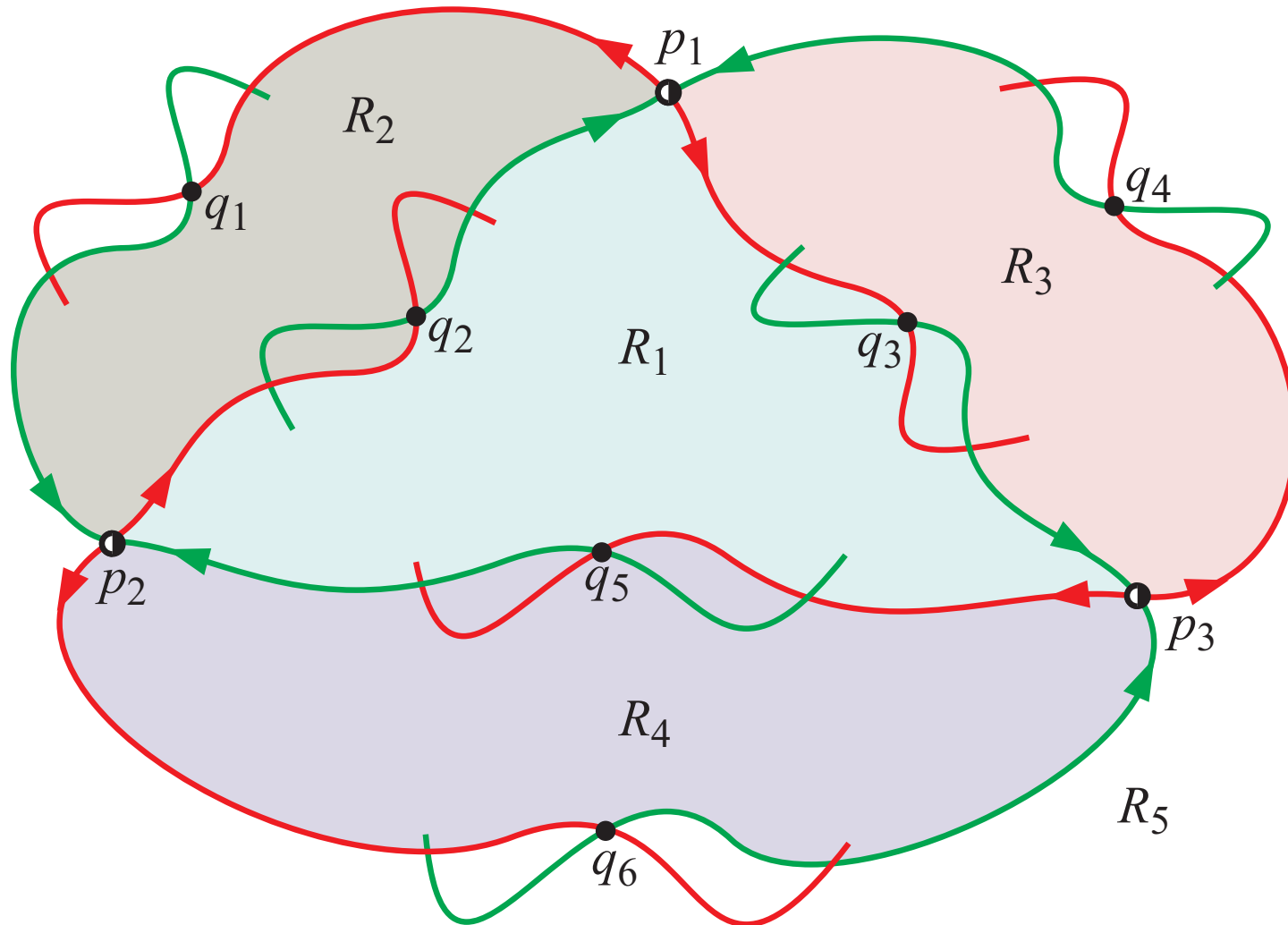
Partition phase space into coherent regions

- Intersection of unstable and stable manifolds define **boundaries**.



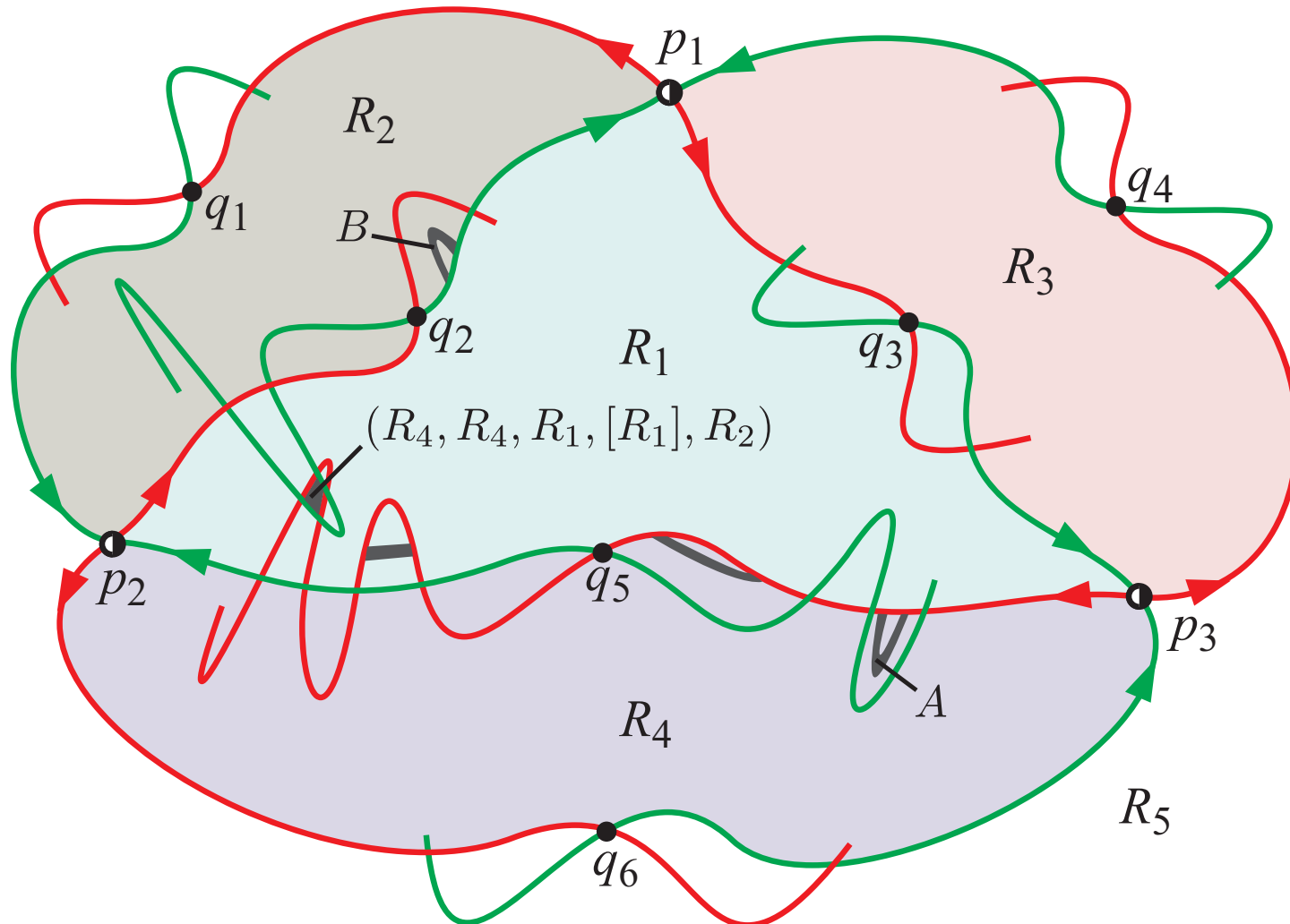
Partition phase space into coherent regions

- These boundaries divide the phase space into **coherent regions**.



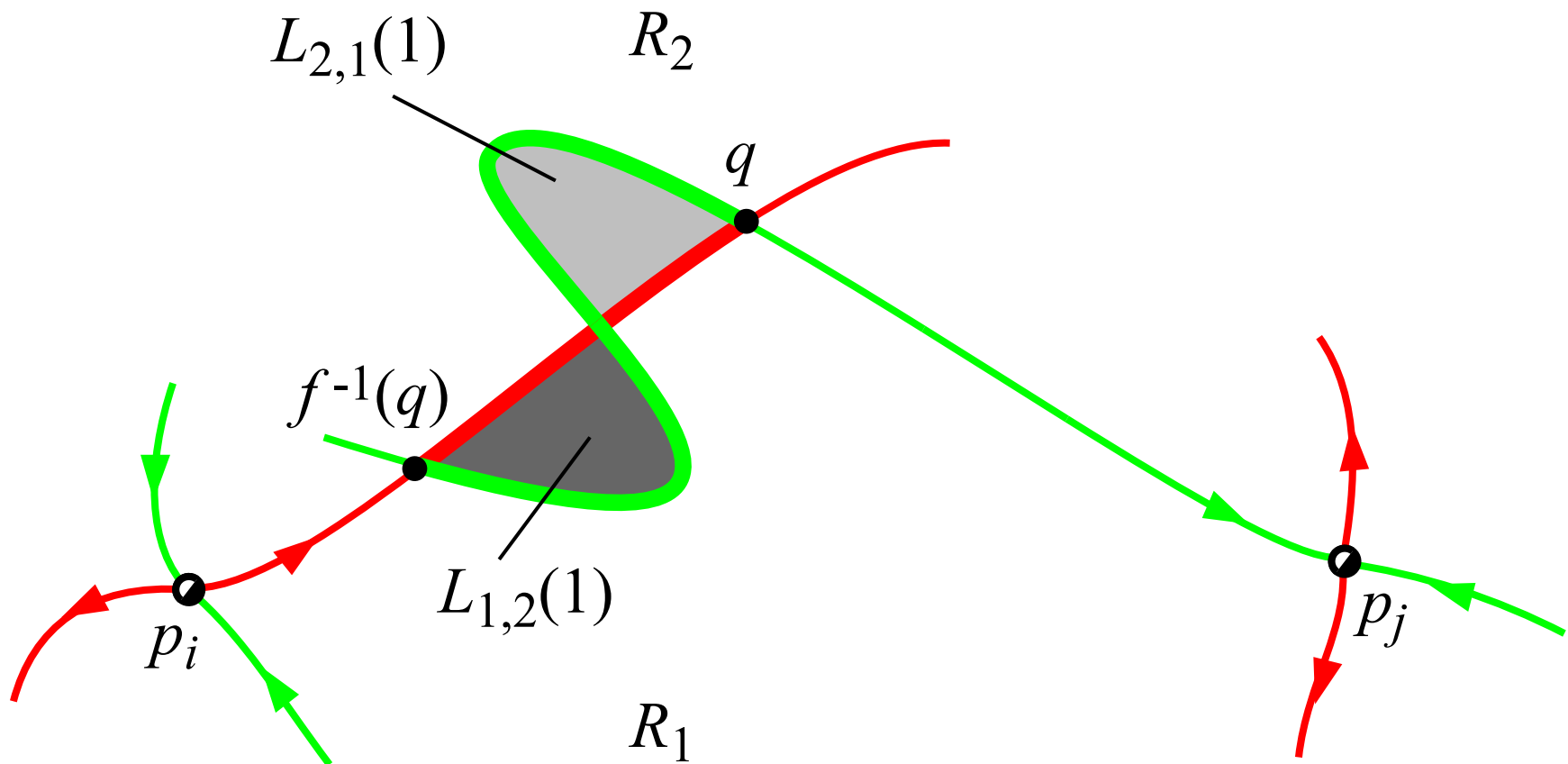
Label mobile subregions: 'atoms' of transport

- Can label mobile subregions based on their past and future whereabouts under one iterate of the map, e.g., $(\dots, R_4, R_4, R_1, [R_1], R_2, \dots)$



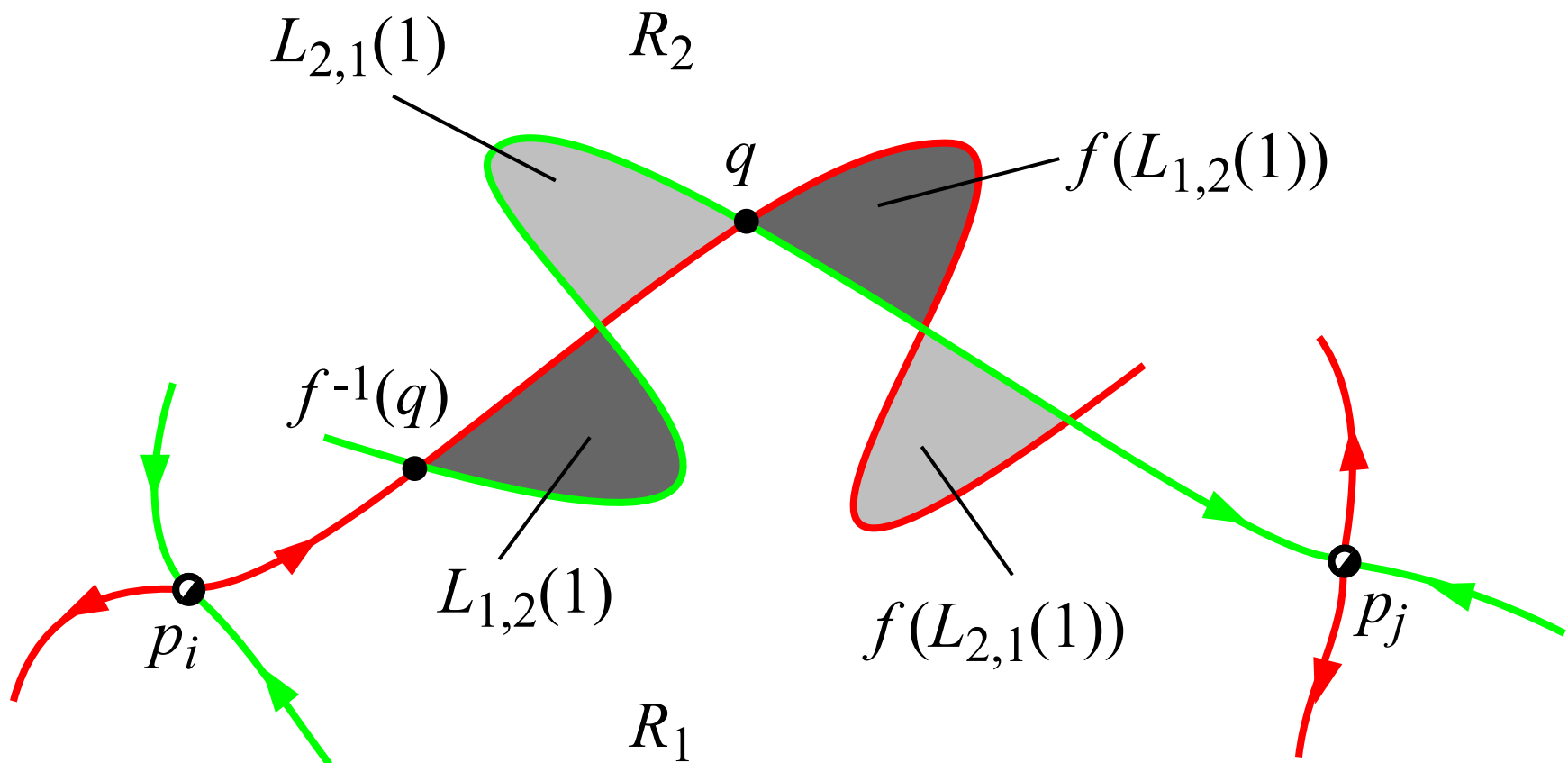
Lobe dynamics: transport across a boundary

- $U[f^{-1}(q), q] \cup S[f^{-1}(q), q]$ forms boundary of two lobes; one in R_1 , labeled $L_{1,2}(1)$, or equivalently $([R_1], R_2)$, where $f(([R_1], R_2)) = (R_1, [R_2])$, etc. for $L_{2,1}(1)$



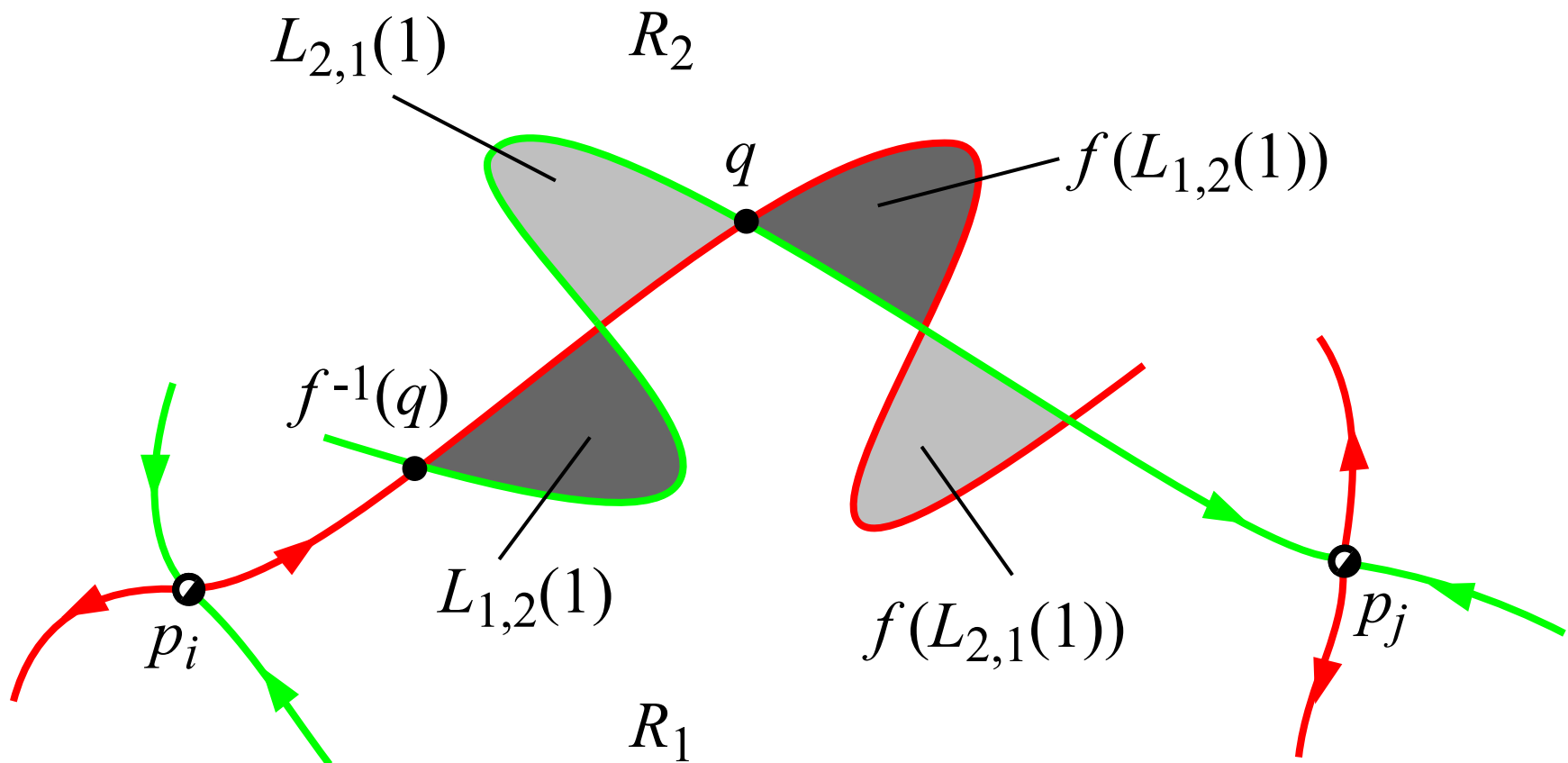
Lobe dynamics: transport across a boundary

- Under one iteration of f , *only points in $L_{1,2}(1)$ can move from R_1 into R_2 by crossing B , etc.*
- The two lobes $L_{1,2}(1)$ and $L_{2,1}(1)$ are called a **turnstile**.



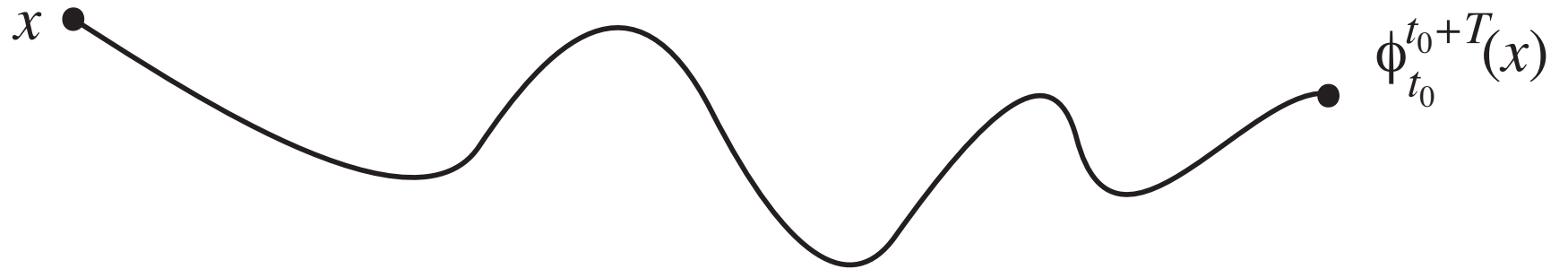
Lobe dynamics: transport across a boundary

- Essence of lobe dynamics: **dynamics associated with crossing a boundary is reduced to the dynamics of turnstile lobes associated with the boundary.**



Extending to realistic flows

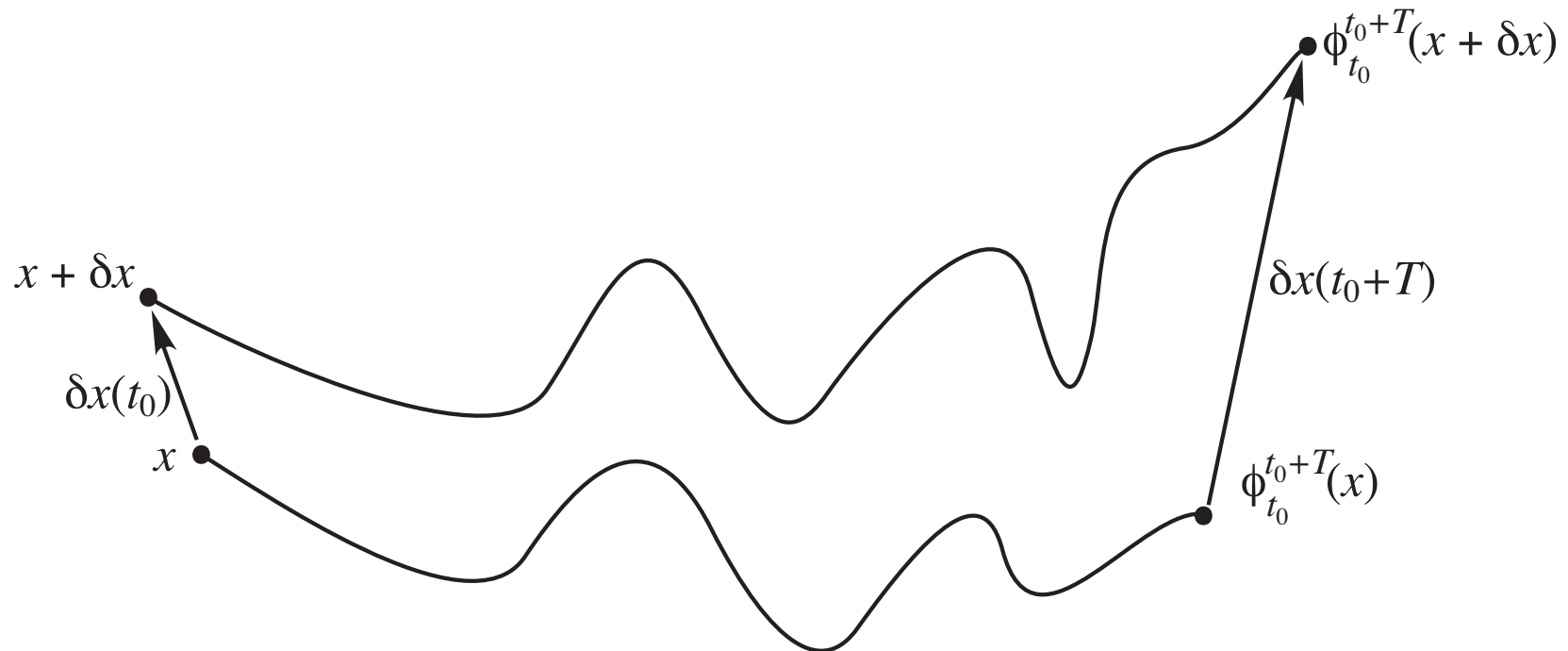
- Data-driven, finite-time, aperiodic setting
- How do we get at transport?
- Recall the flow, $x \mapsto \phi_t^{t+T}(x)$



Identify regions of high sensitivity of initial conditions

- Small initial perturbations $\delta x(t)$ grow like

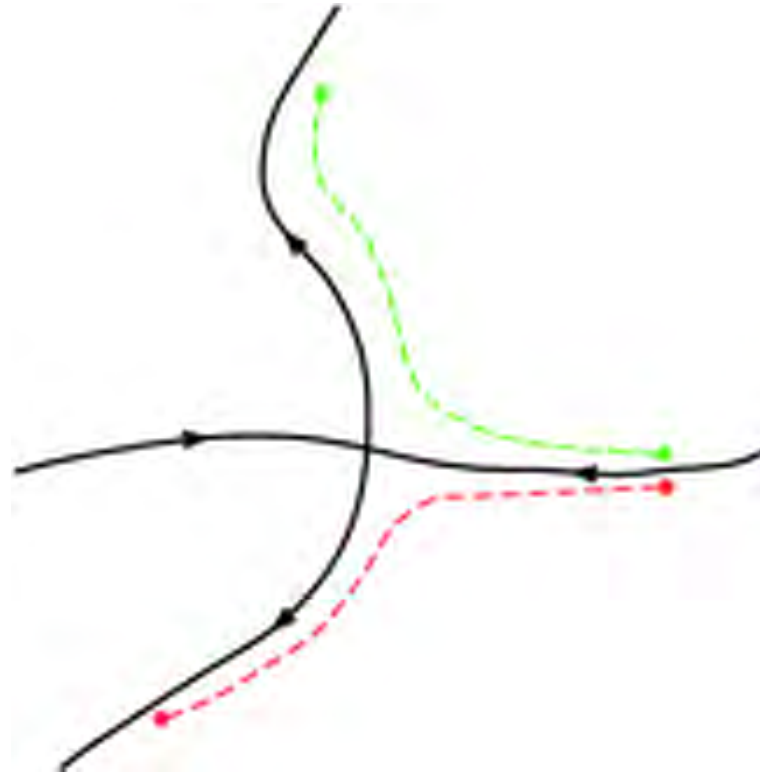
$$\begin{aligned}\delta x(t + T) &= \phi_t^{t+T}(x + \delta x(t)) - \phi_t^{t+T}(x) \\ &= \frac{d\phi_t^{t+T}(x)}{dx} \delta x(t) + O(\|\delta x(t)\|^2)\end{aligned}$$



Identify regions of high sensitivity of initial conditions

- Small initial perturbations $\delta x(t)$ grow like

$$\begin{aligned}\delta x(t + T) &= \phi_t^{t+T}(x + \delta x(t)) - \phi_t^{t+T}(x) \\ &= \frac{d\phi_t^{t+T}(x)}{dx} \delta x(t) + O(\|\delta x(t)\|^2)\end{aligned}$$



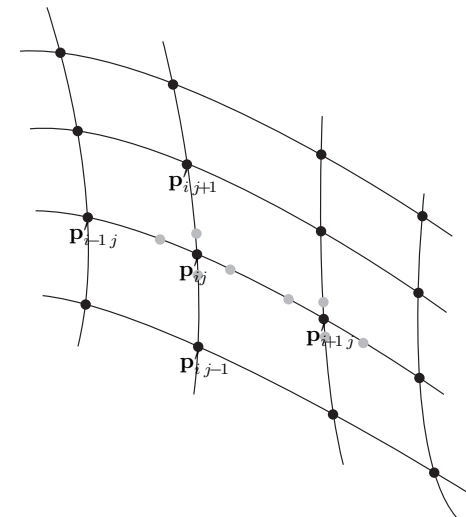
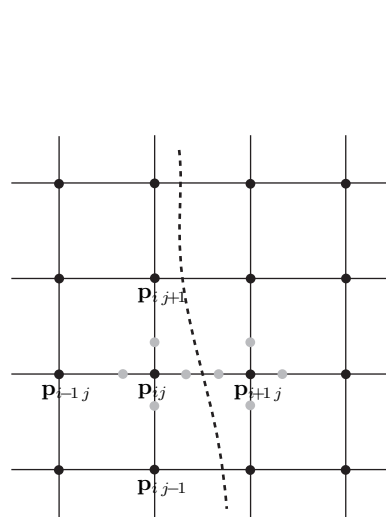
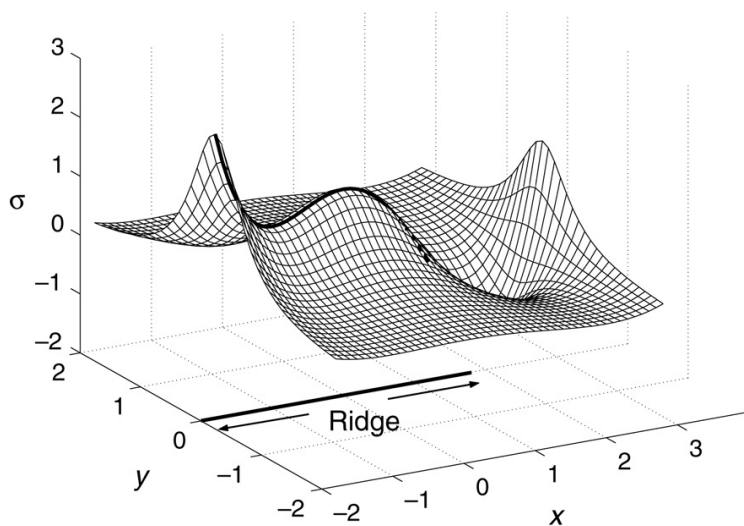
Invariant manifold analogs: FTLE-LCS approach

- The finite-time Lyapunov exponent (FTLE),

$$\sigma_t^T(x) = \frac{1}{|T|} \log \left\| \frac{d\phi_t^{t+T}(x)}{dx} \right\|$$

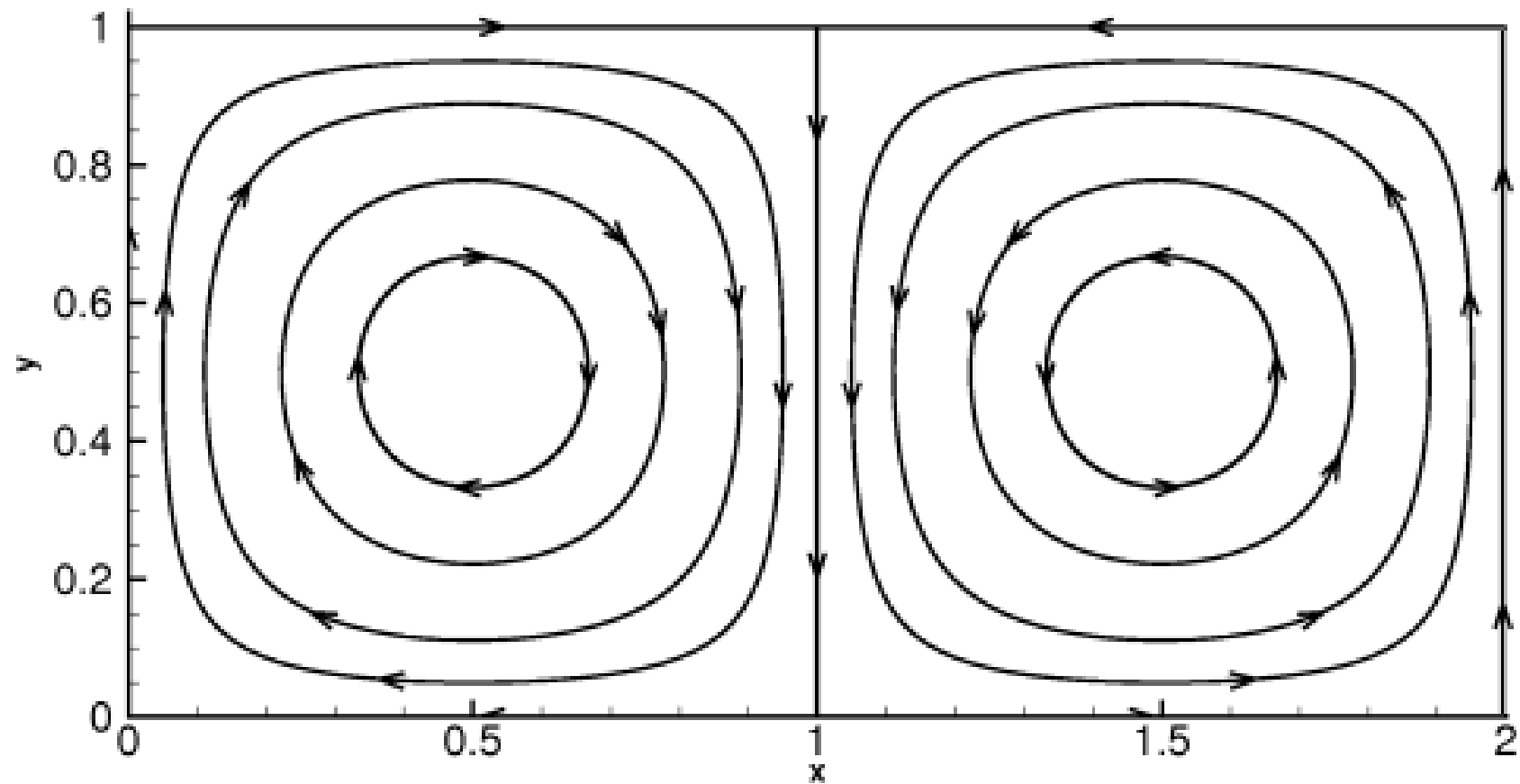
measures the maximum stretching rate over the interval T of trajectories starting near the point x at time t

- Ridges of σ_t^T are candidate hyperbolic codim-1 surfaces; finite-time analogs of stable/unstable manifolds; Lagrangian coherent structures¹

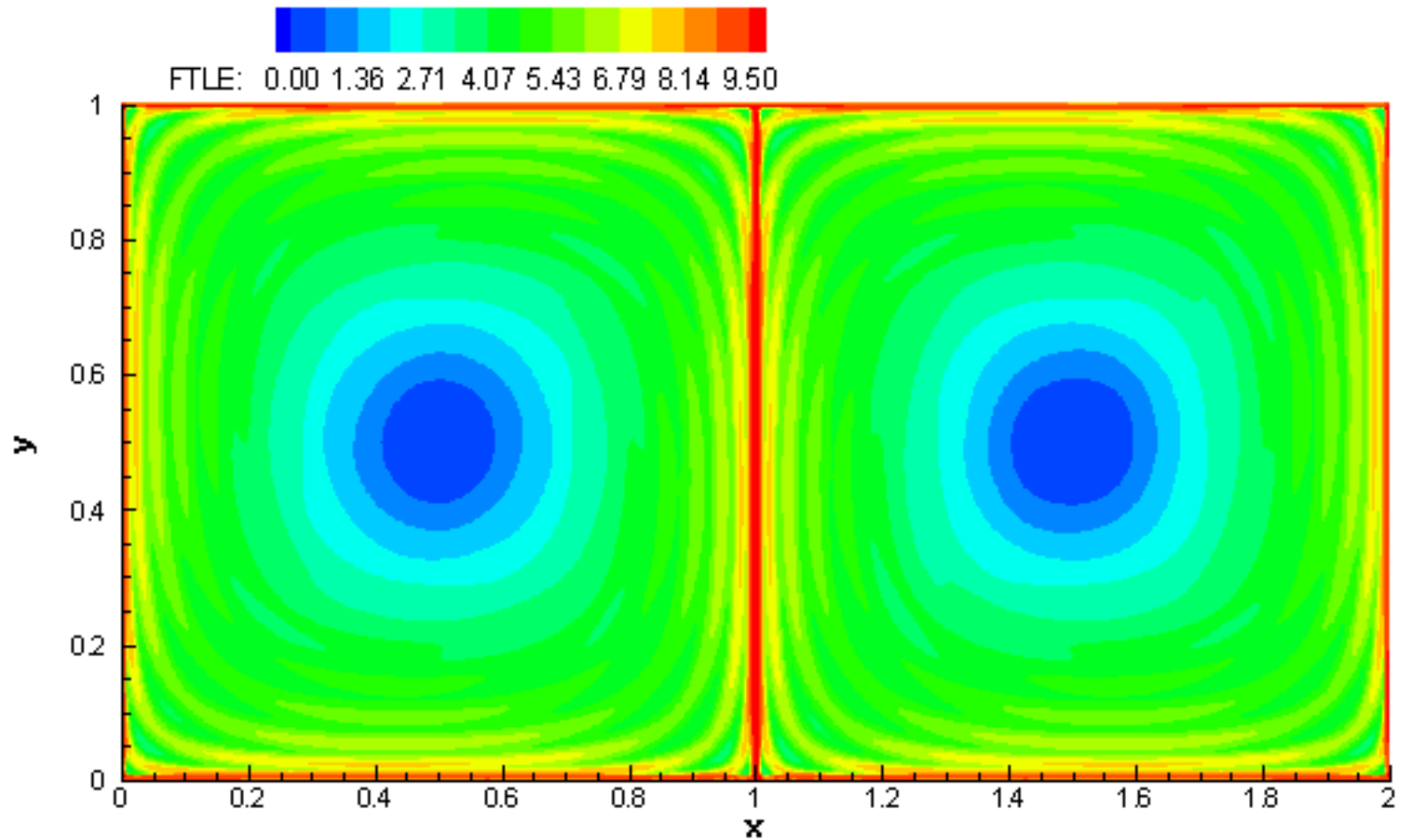


¹cf. Bowman, 1999; Haller & Yuan, 2000; Haller, 2001; Shadden, Lekien, Marsden, 2005

Invariant manifold analogs: FTLE-LCS approach



Invariant manifold analogs: FTLE-LCS approach

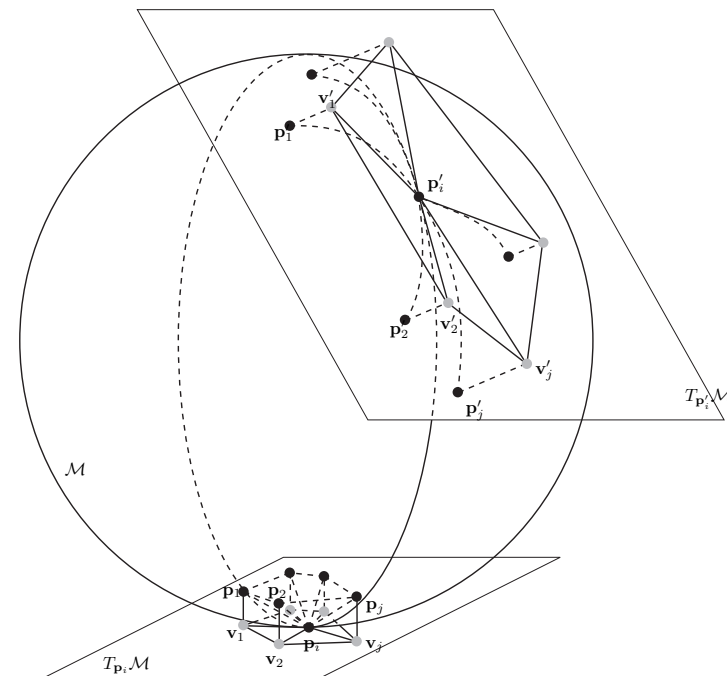
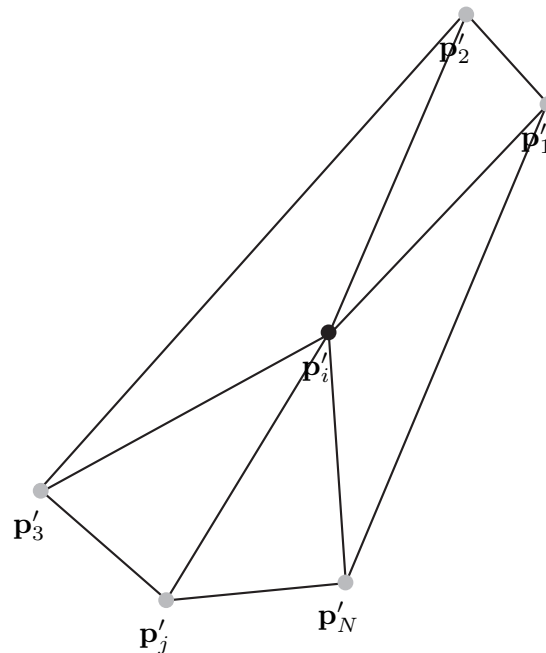
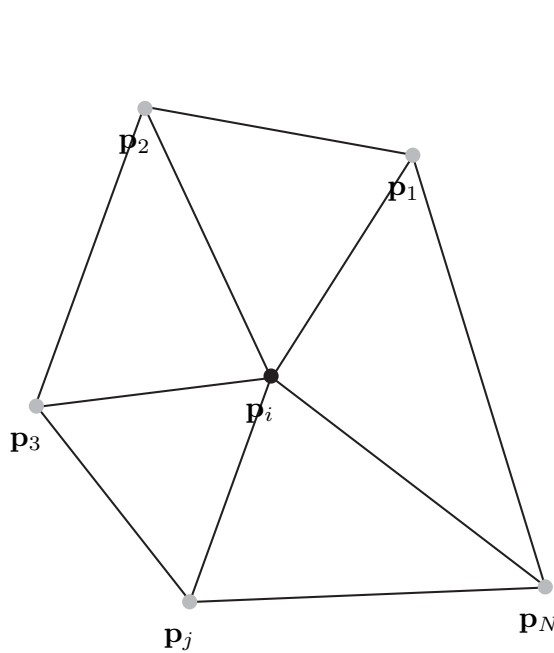


Invariant manifold analogs: FTLE-LCS approach

- We can define the FTLE for Riemannian manifolds²

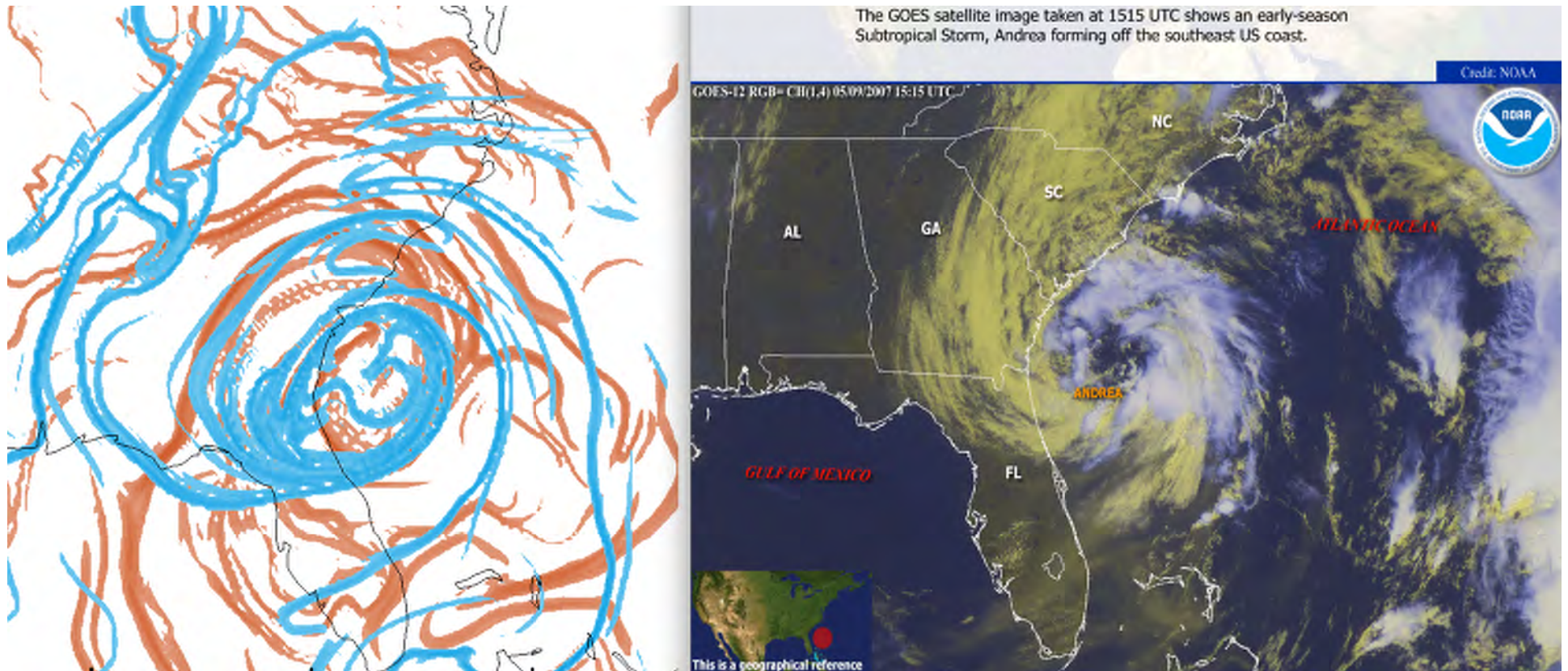
$$\sigma_t^T(x) = \frac{1}{|T|} \ln \left\| D\phi_t^{t+T} \right\| \doteq \frac{1}{|T|} \ln \left(\max_{y \neq 0} \frac{\left\| D\phi_t^{t+T}(y) \right\|}{\|y\|} \right)$$

with y a small perturbation in the tangent space at x .



²Lekien & Ross [2010] Chaos

Hurricanes and lobe dynamics



orange = repelling LCSs, blue = attracting LCSs

satellite

Andrea, first storm of 2007 hurricane season

cf. Sapsis & Haller [2009], Du Toit & Marsden [2010], Lekien & Ross [2010], Tallapragada & Ross [2011]

Hurricanes and lobe dynamics



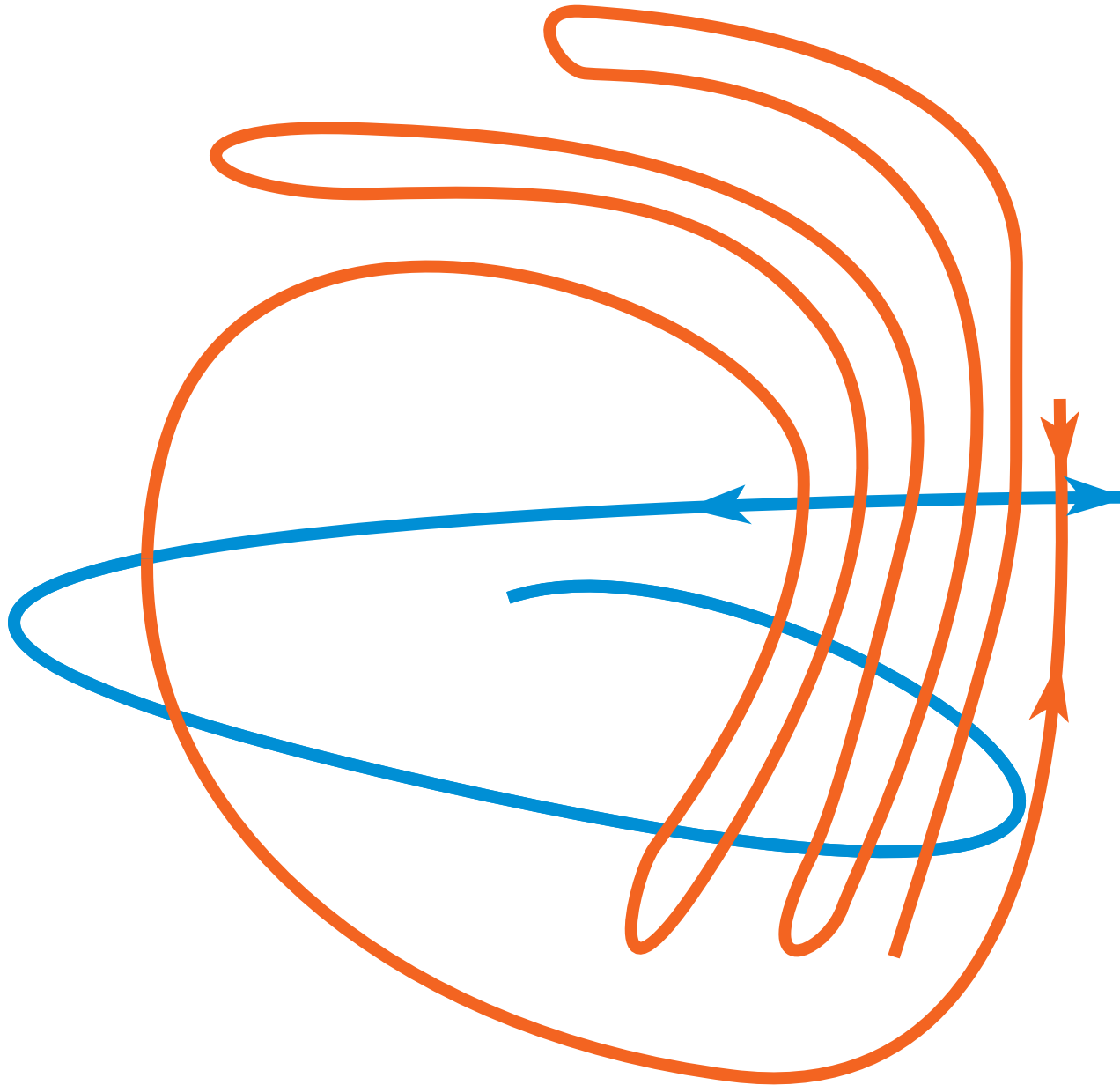
Andrea at one snapshot; LCS shown (orange = repelling, blue = attracting)

Hurricanes and lobe dynamics



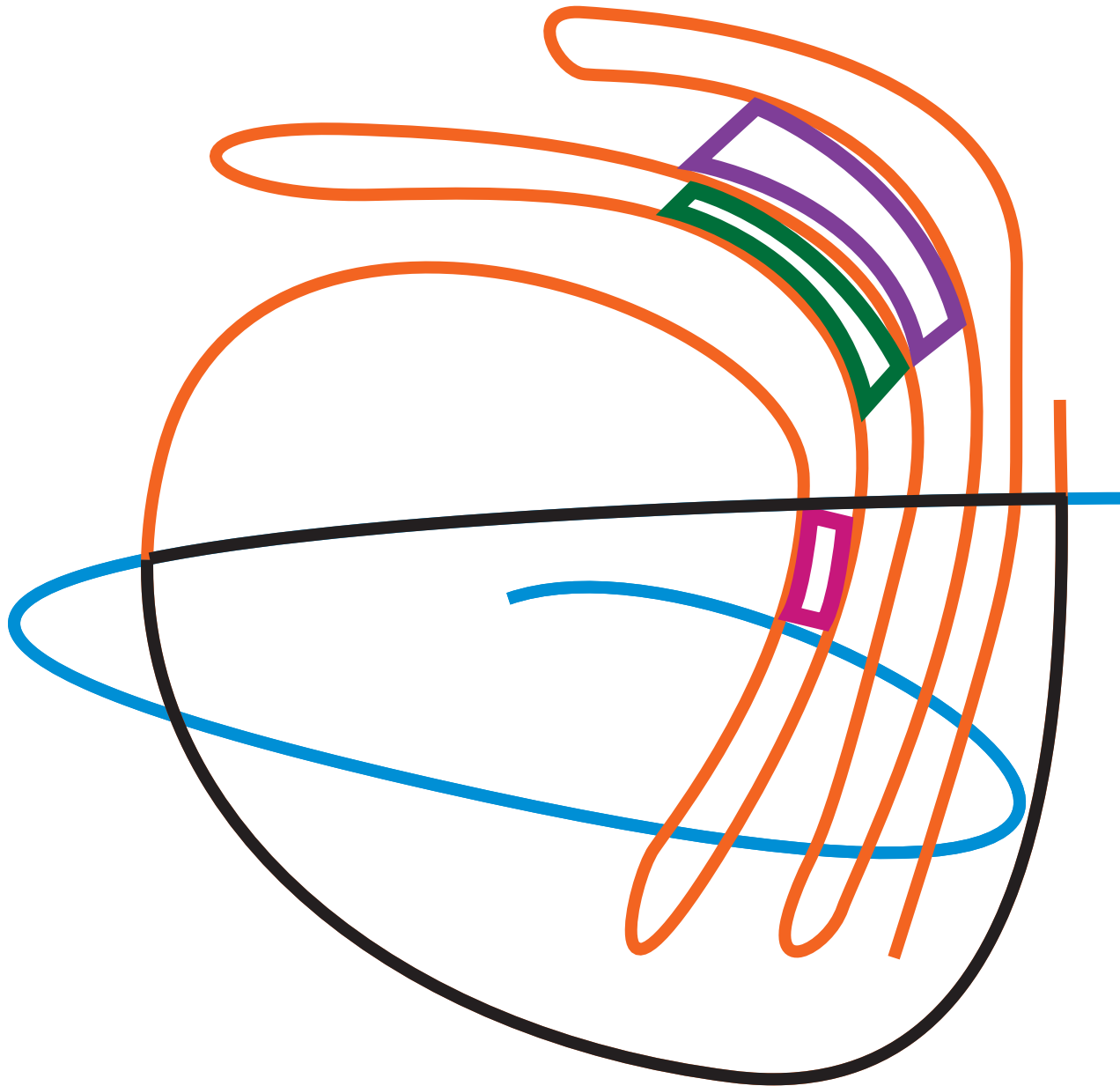
orange = repelling (stable manifold), blue = attracting (unstable manifold)

Hurricanes and lobe dynamics



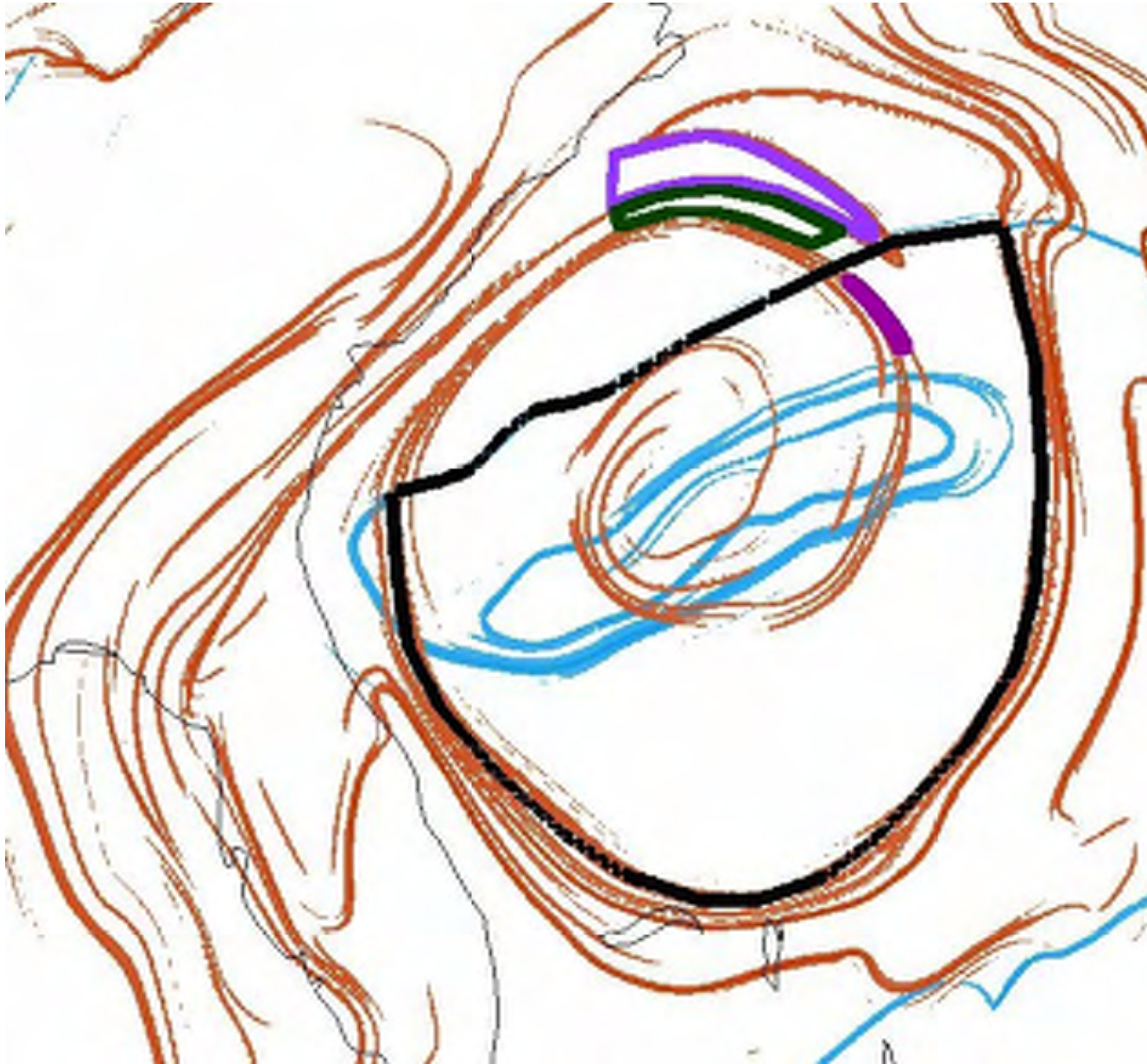
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Hurricanes and lobe dynamics



Portions of lobes colored; magenta = outgoing, green = incoming, purple = stays out

Hurricanes and lobe dynamics

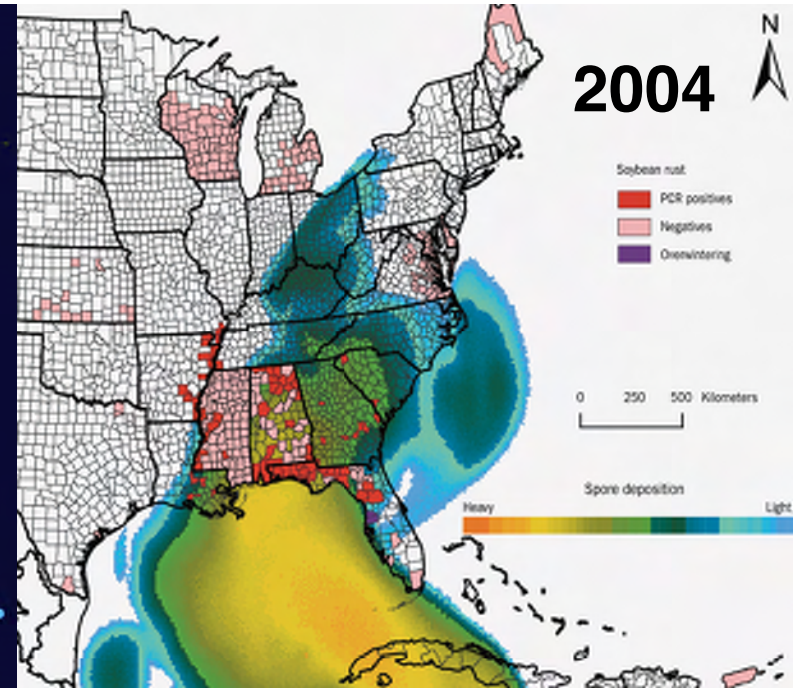


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Hurricanes and lobe dynamics

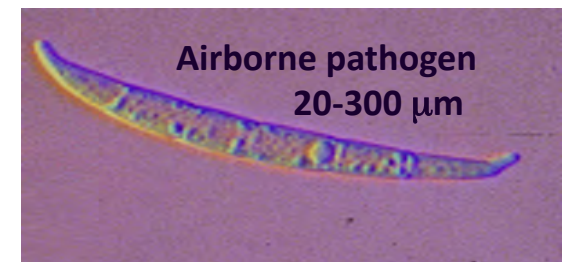
Sets behave as lobe dynamics dictates

Invasive species riding coherent structures

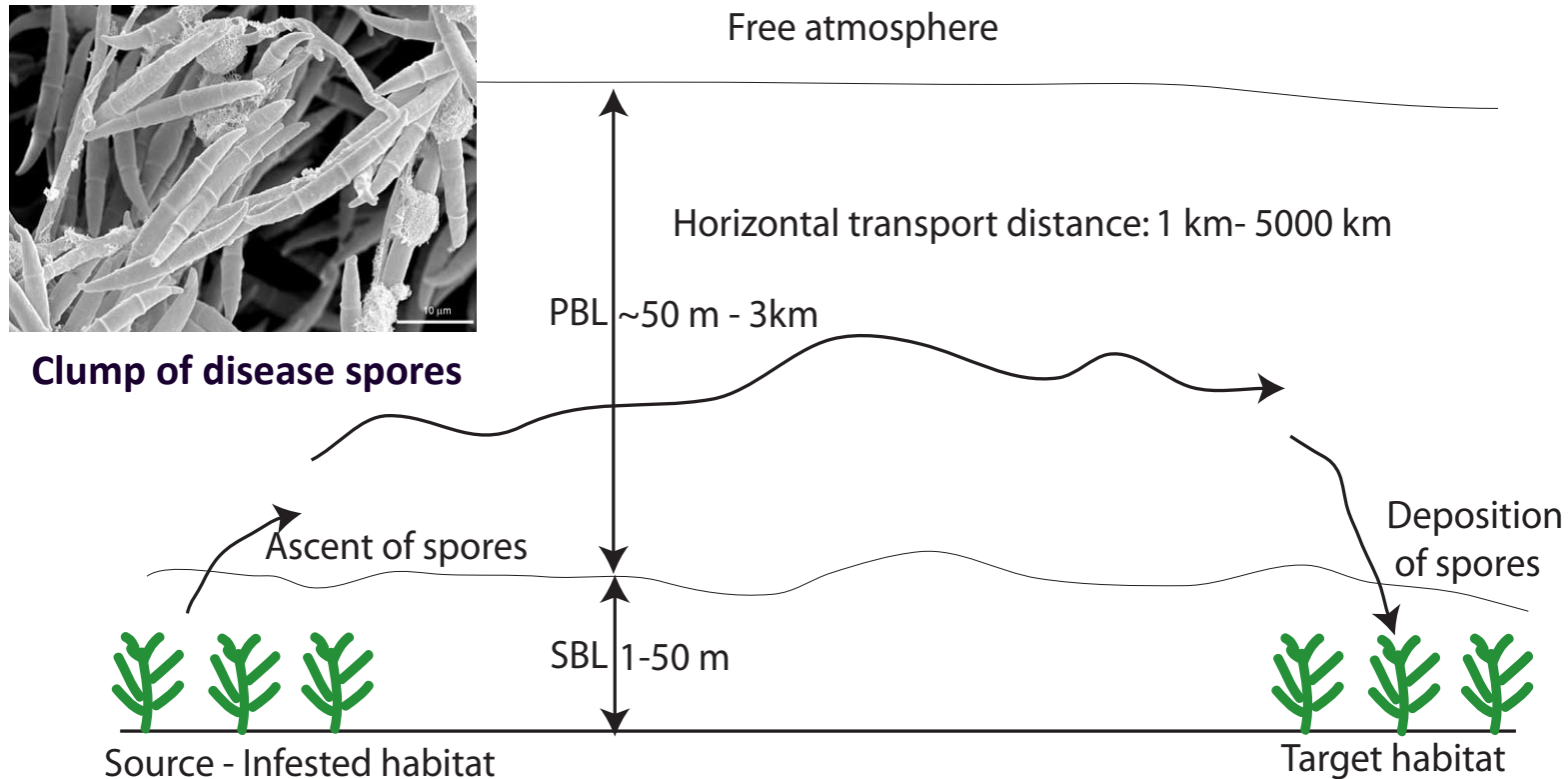


Disease extent

Cost of invasive organisms is **\$137 billion** per year in U.S.



Atmospheric transport of microorganisms



e.g., *Fusarium*

- Spore production, release, escape from surface
- **Long-range transport (time-scale hours to days)**
- Deposition, infection efficiency, host susceptibility

Atmospheric transport network relevant for aeroecology

Skeleton of large-scale
horizontal transport

relevant for large-scale
spatiotemporal patterns
of important biota
e.g., plant pathogens

orange = repelling LCSs, blue = attracting LCSs

Aerial sampling:
40 m – 400 m altitude



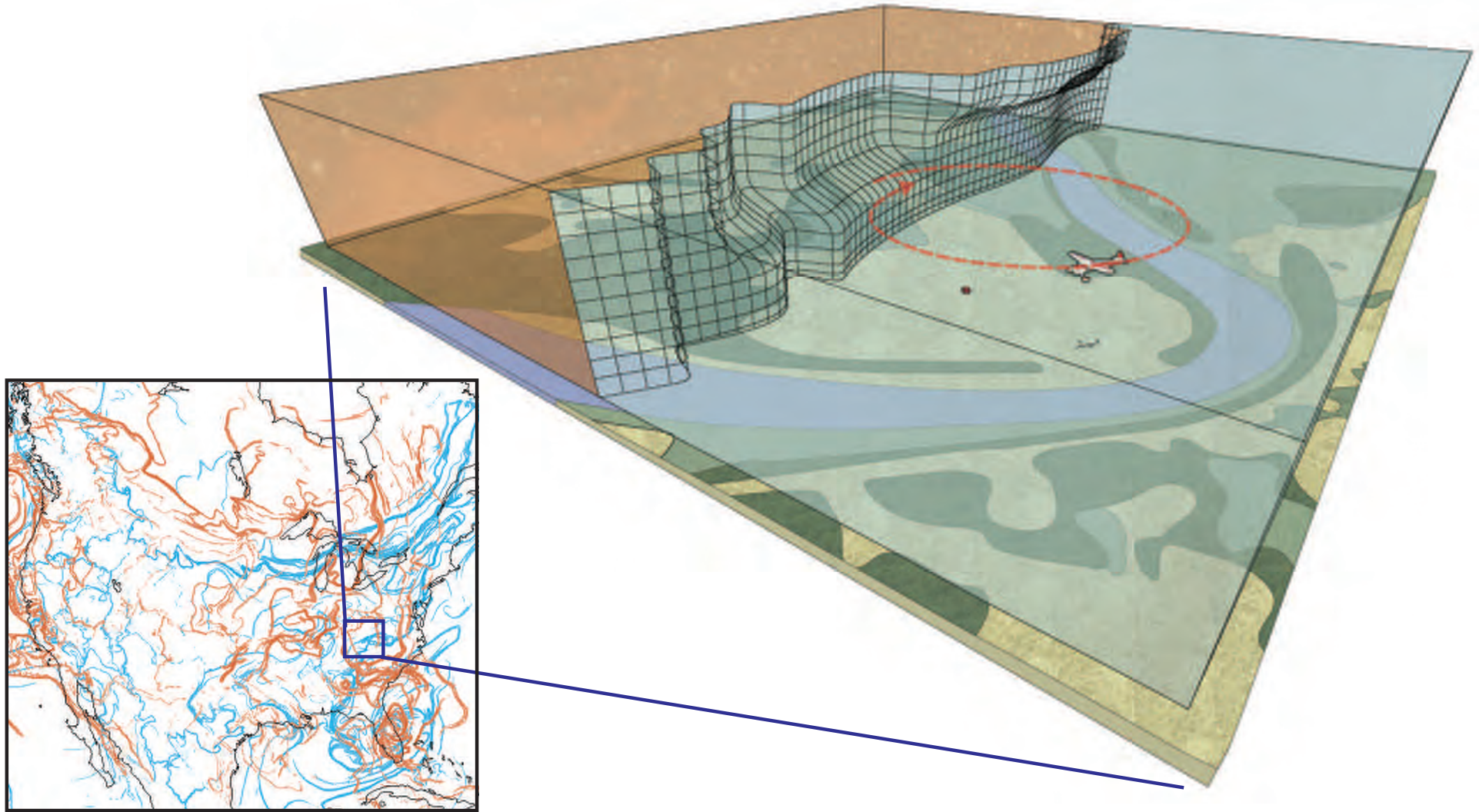
Image © 2010 Commonwealth of Virginia
Image © 2010 DigitalGlobe
Image USDA Farm Service Agency
Image U.S. Geological Survey

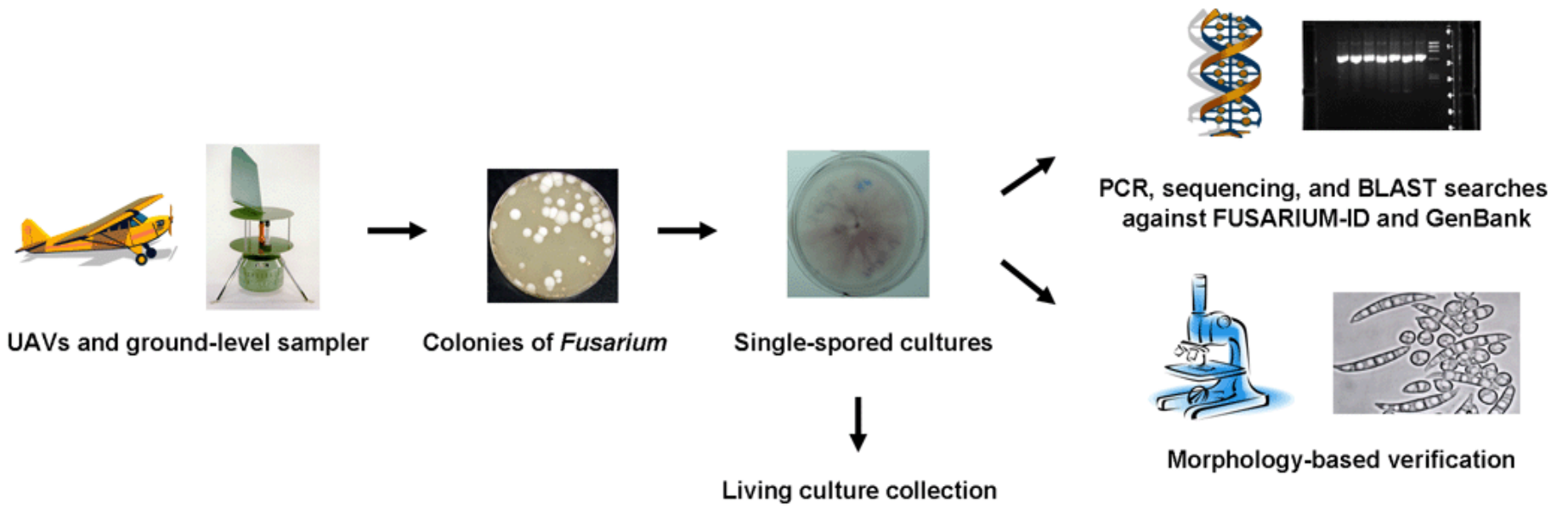
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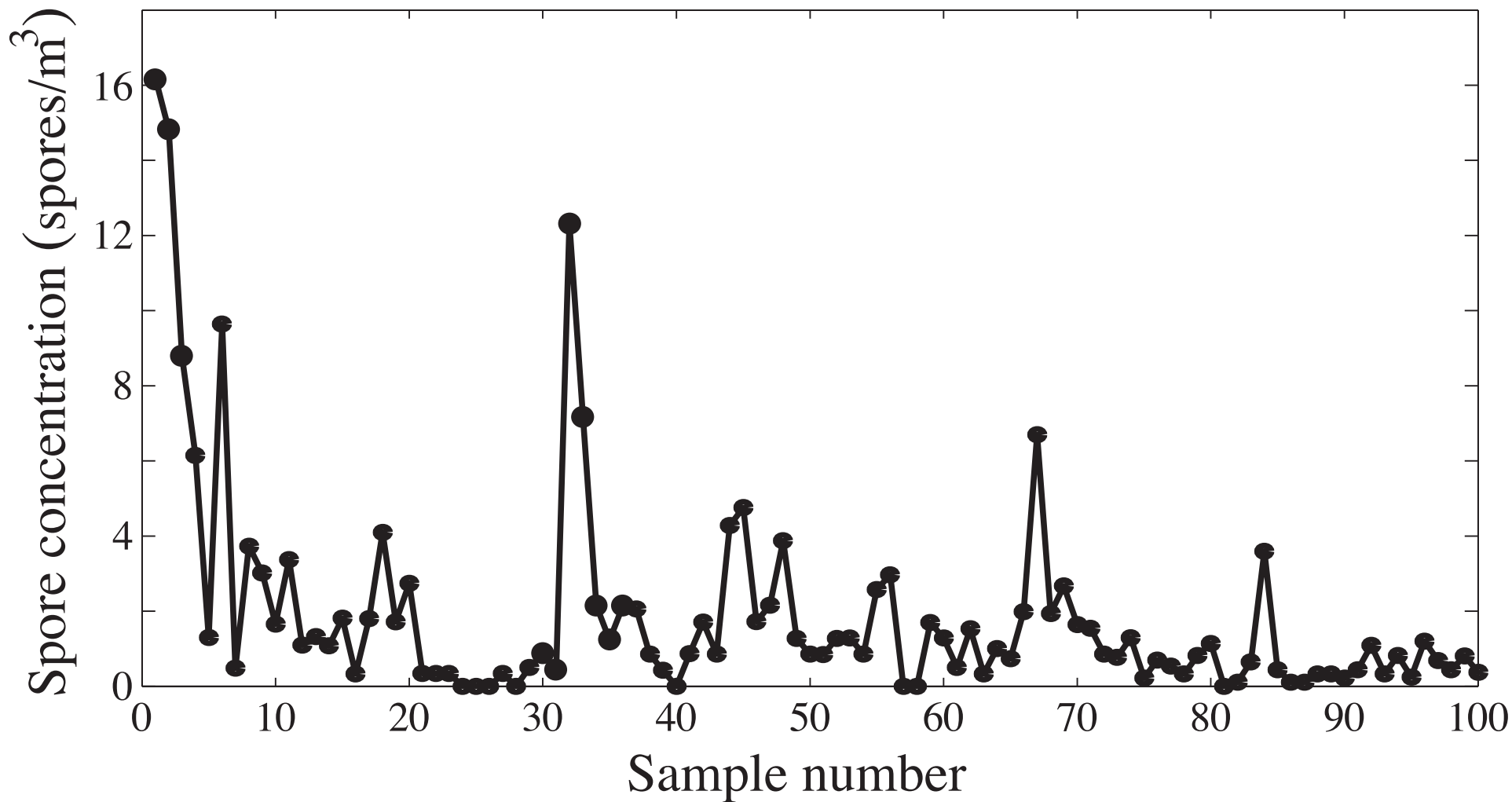
Eye alt 0294 ft



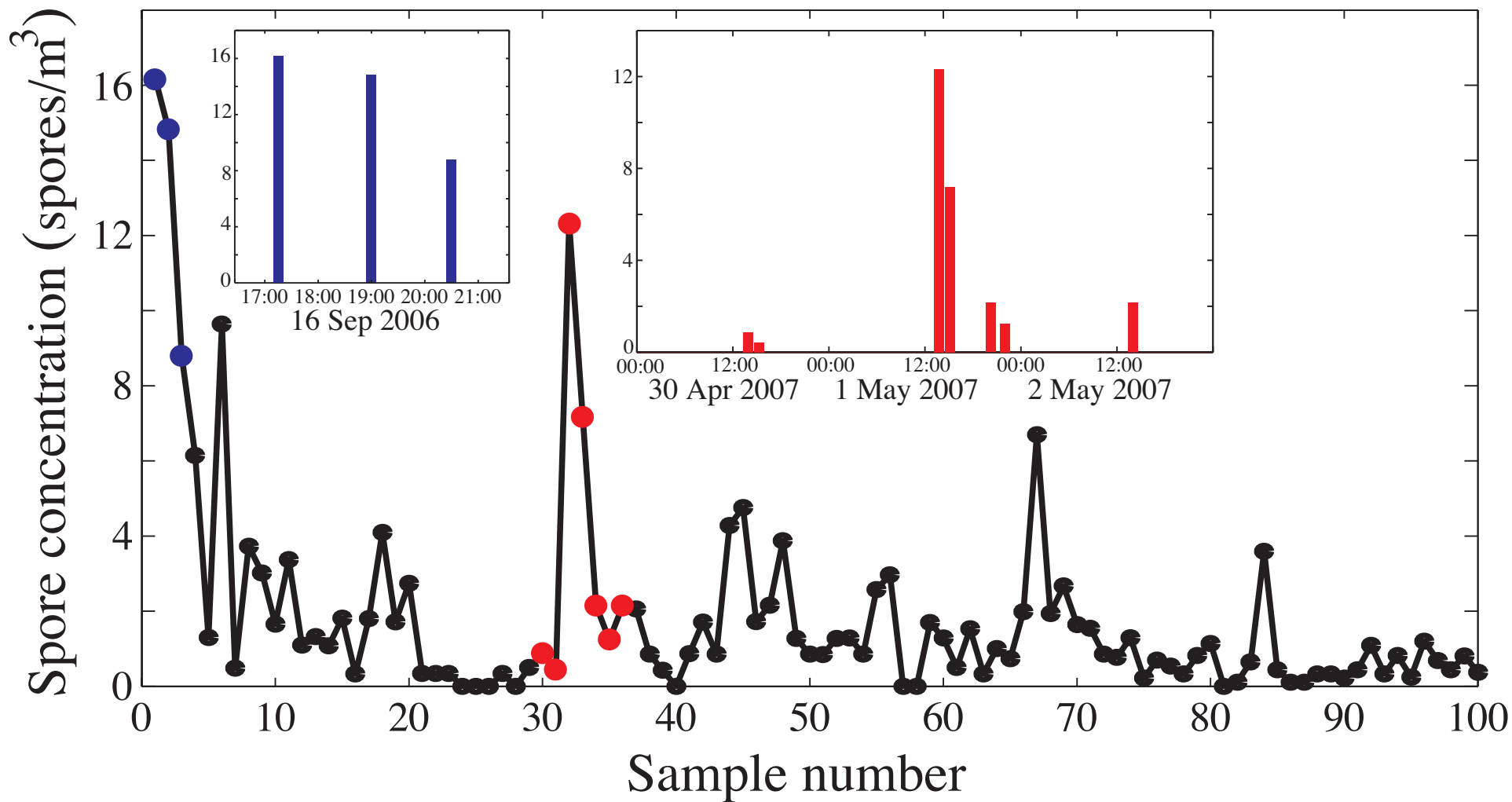
Atmospheric transport network relevant for aeroecology





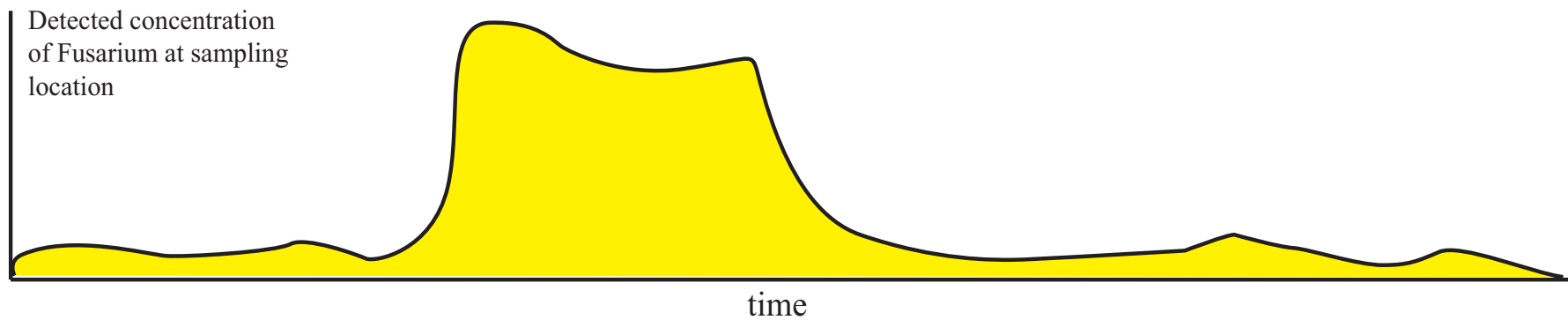


Concentration of *Fusarium* spores (number/m³) for samples from 100 flights conducted between August 2006 and March 2010.

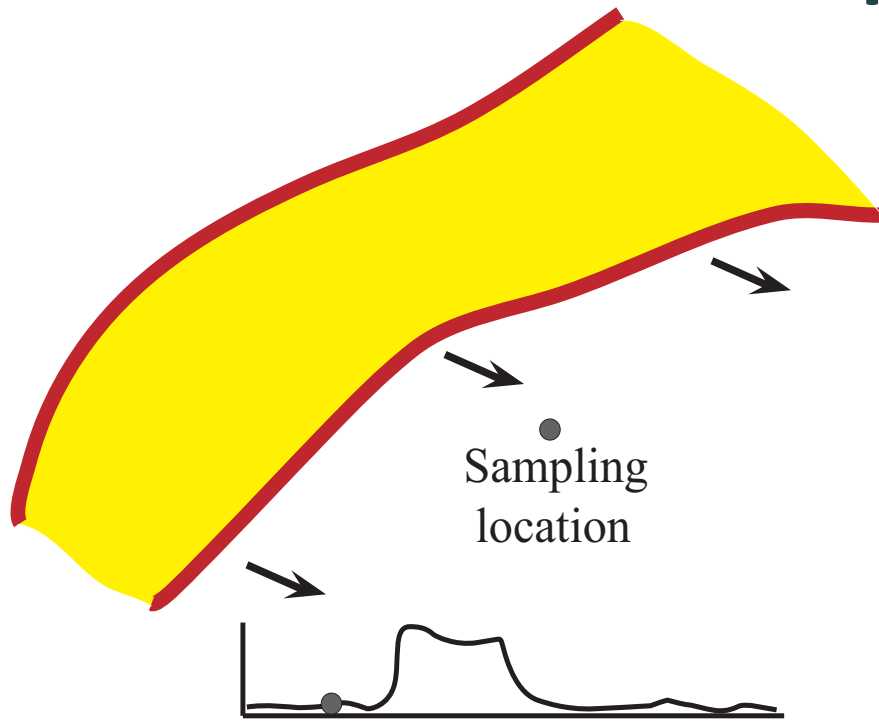


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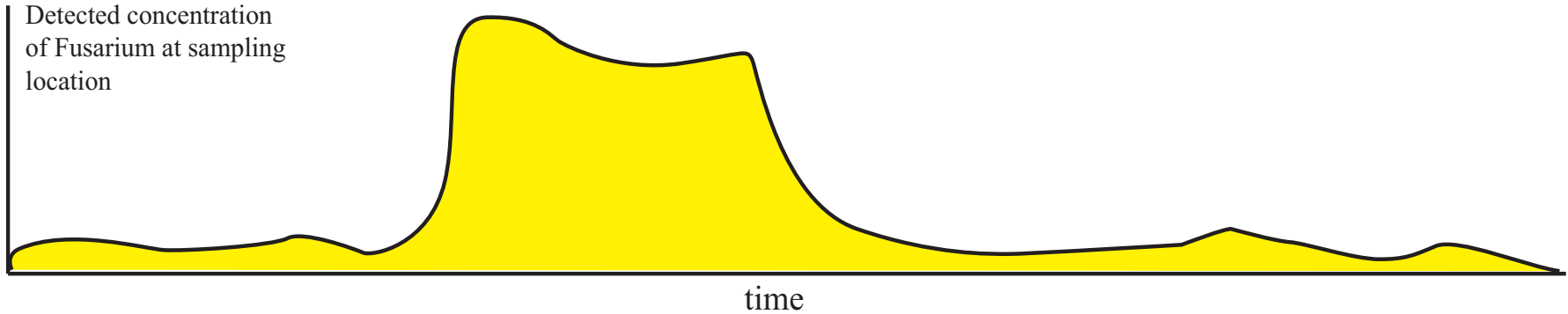
Punctuated changes: correlated to LCS passage?



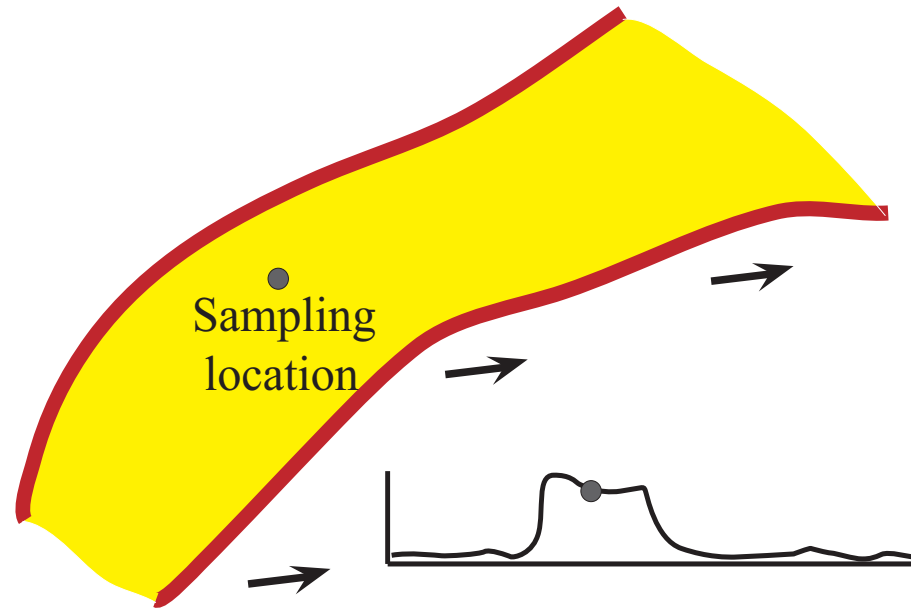
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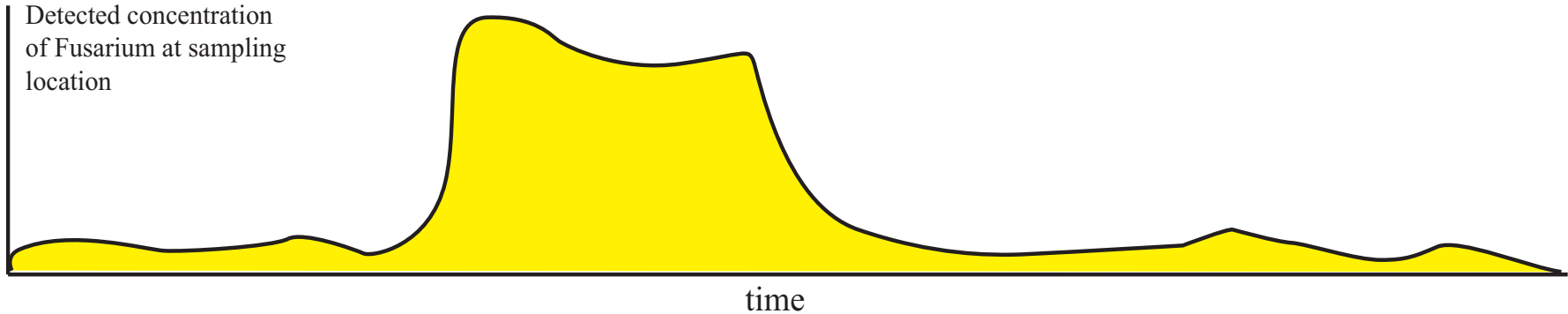
Detected concentration
of Fusarium at sampling
location



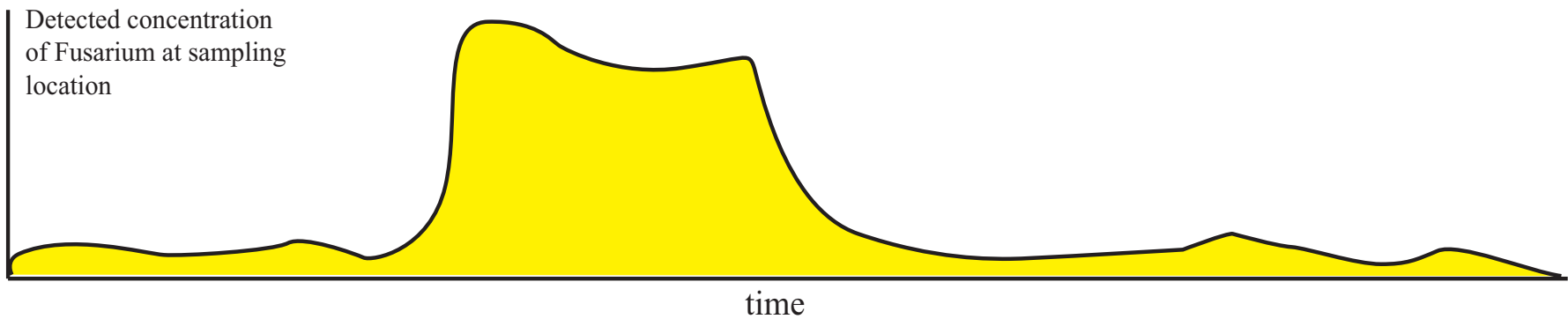
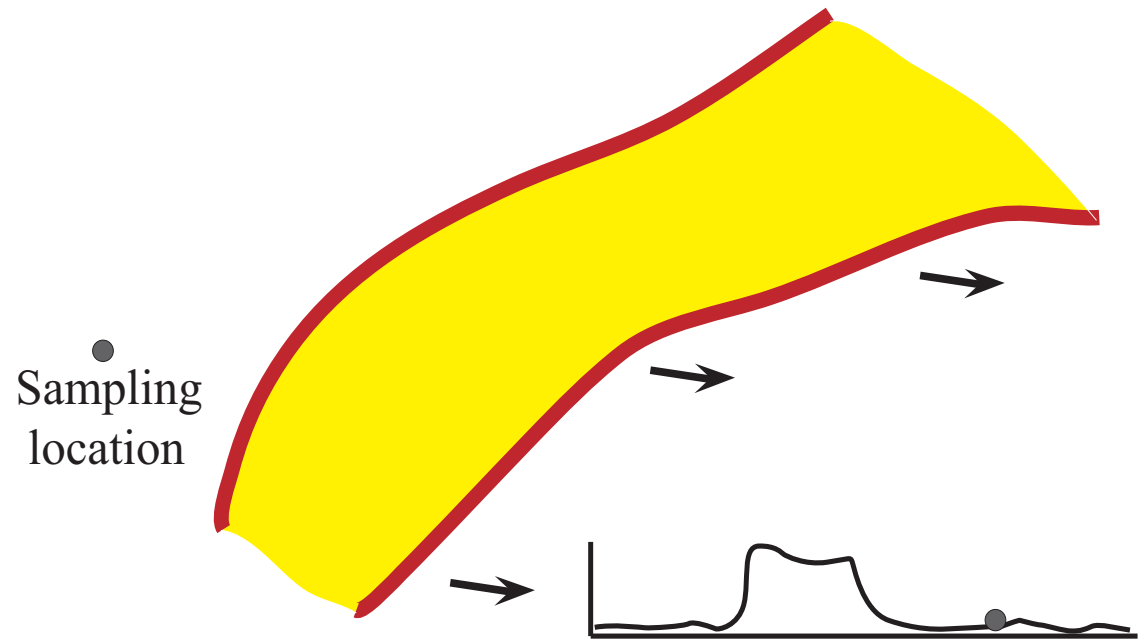
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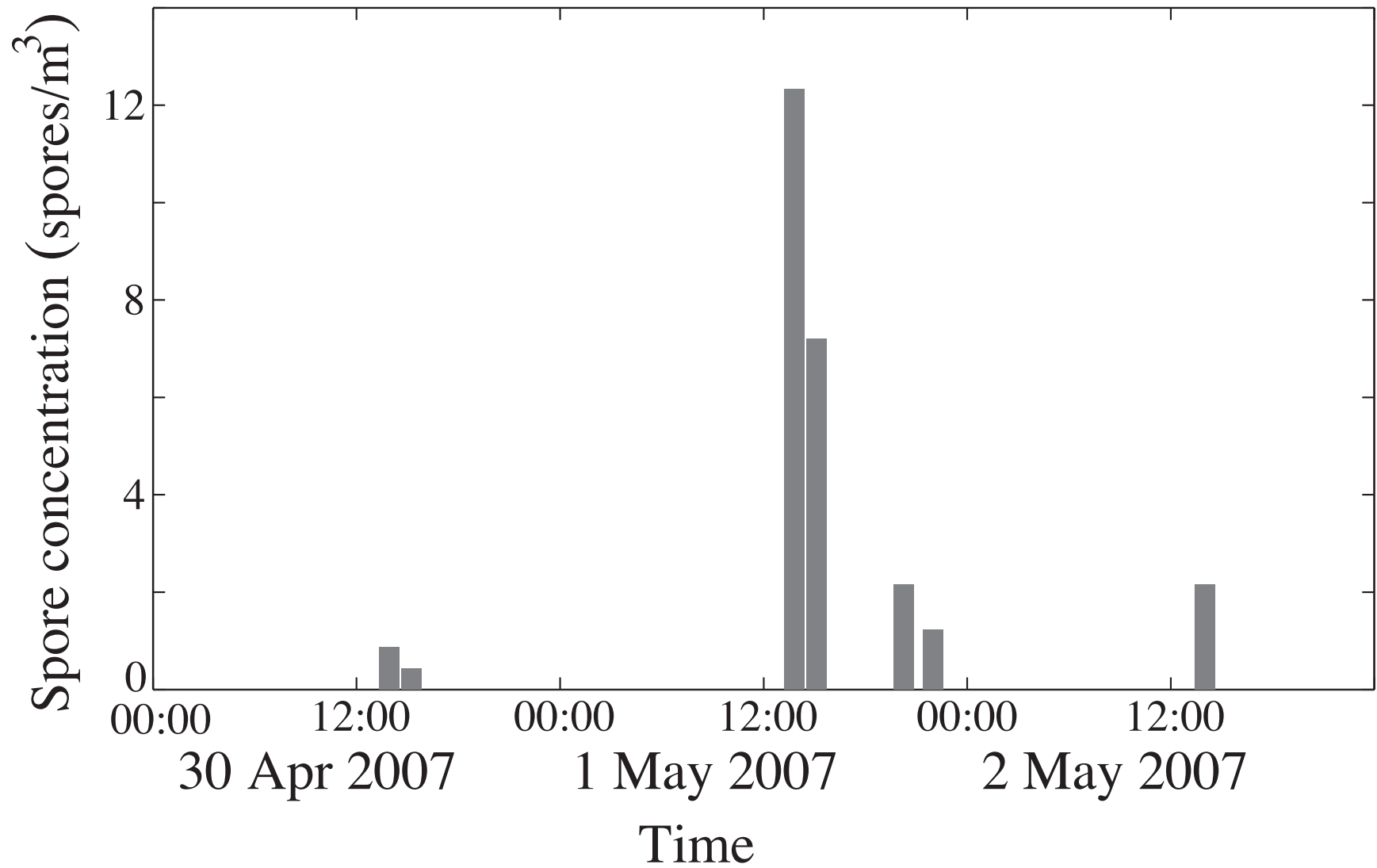


Detected concentration
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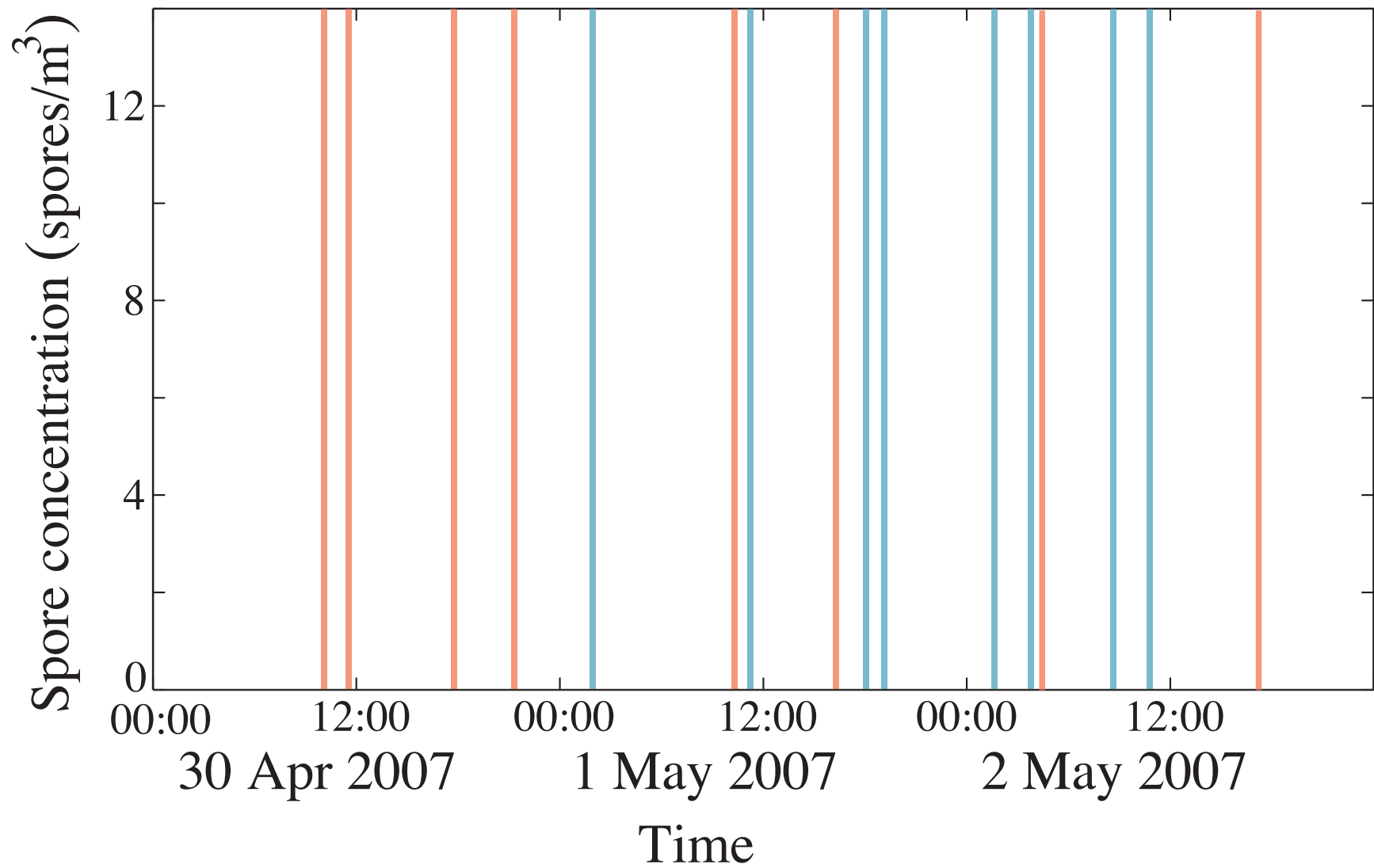


Punctuated changes: correlated to LCS passage?





Time series of concentration $\{(t_0, C_0), \dots, (t_{N-1}, C_{N-1})\}$

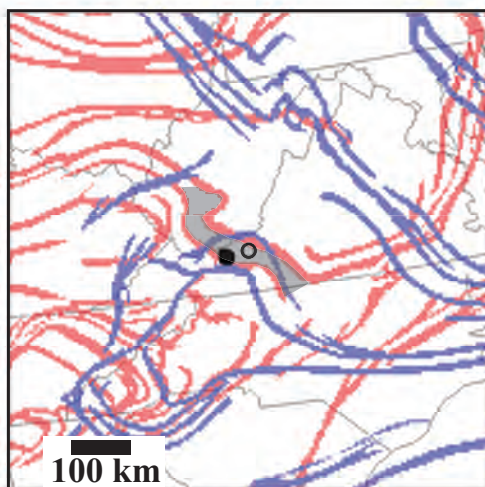


LCS passage times: orange = repelling LCSs, blue = attracting

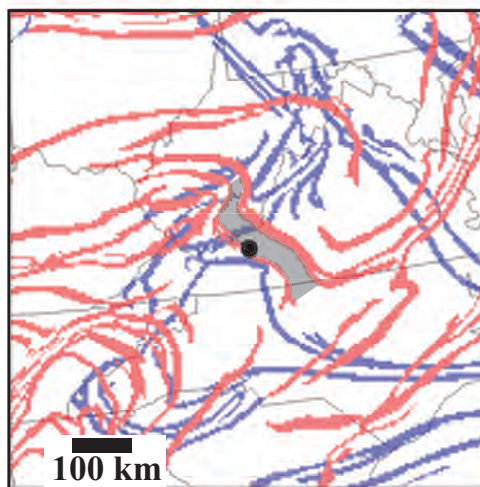
Summary of Hypothesis Testing

- Of 100 samples, only 73 sample pairs within 24 hours
- Of those, 16 show punctuated changes in the concentration of *Fusarium*
- Punctuated change \Rightarrow repelling LCS passage **70% of the time**
($p = 0.0017$)
- **Punctuated changes were significantly associated with the movement of a repelling LCS**
- Correlation poor for attracting LCS: punctuated change \Rightarrow attracting LCS passage 37% of the time ($p = 0.33$)

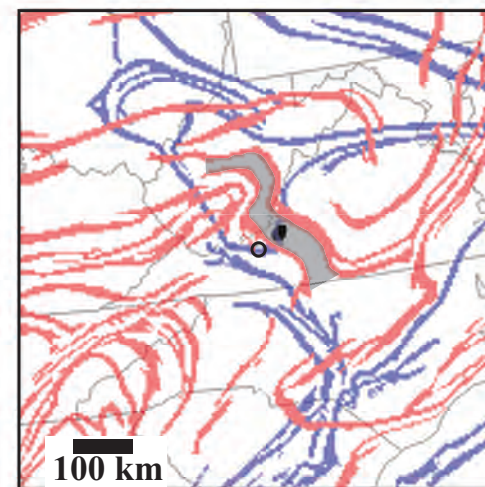
Example: Filament bounded by repelling LCS



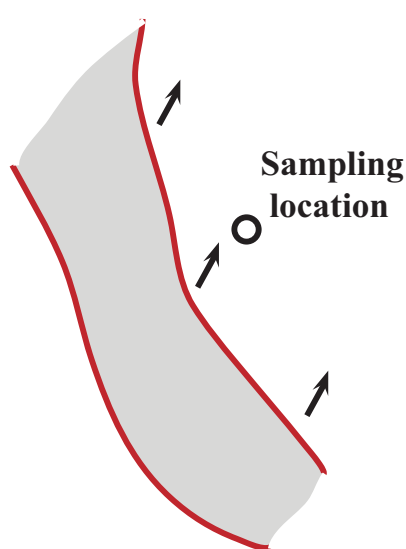
(a)



(b)

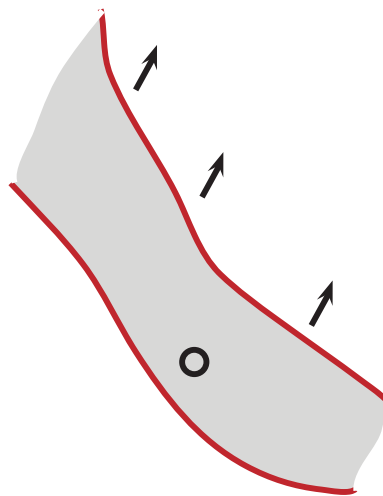


(c)



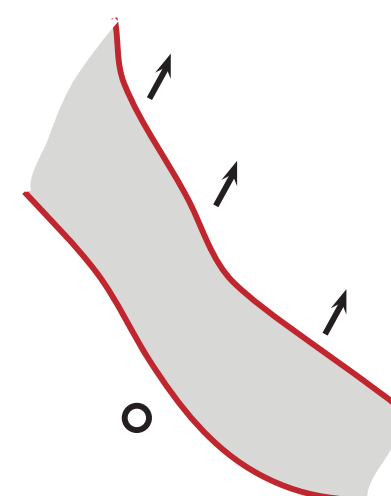
(d)

12:00 UTC 1 May 2007



(e)

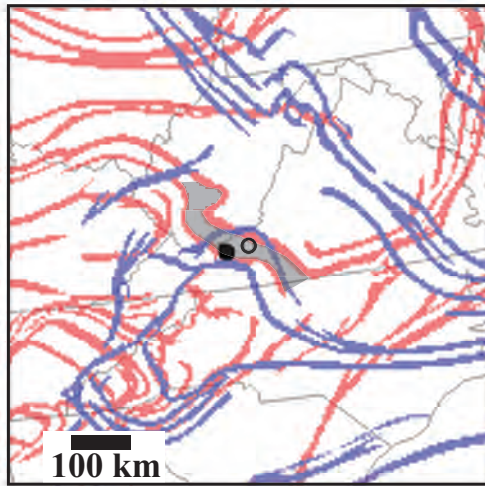
15:00 UTC 1 May 2007



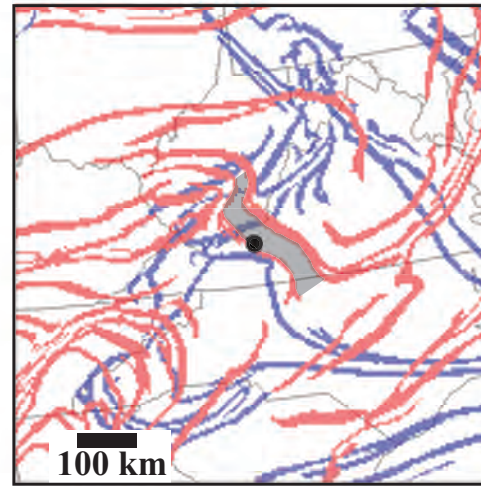
(f)

18:00 UTC 1 May 2007

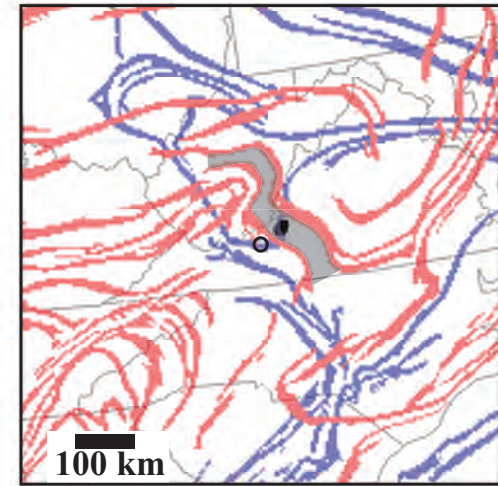
Example: Filament bounded by repelling LCS



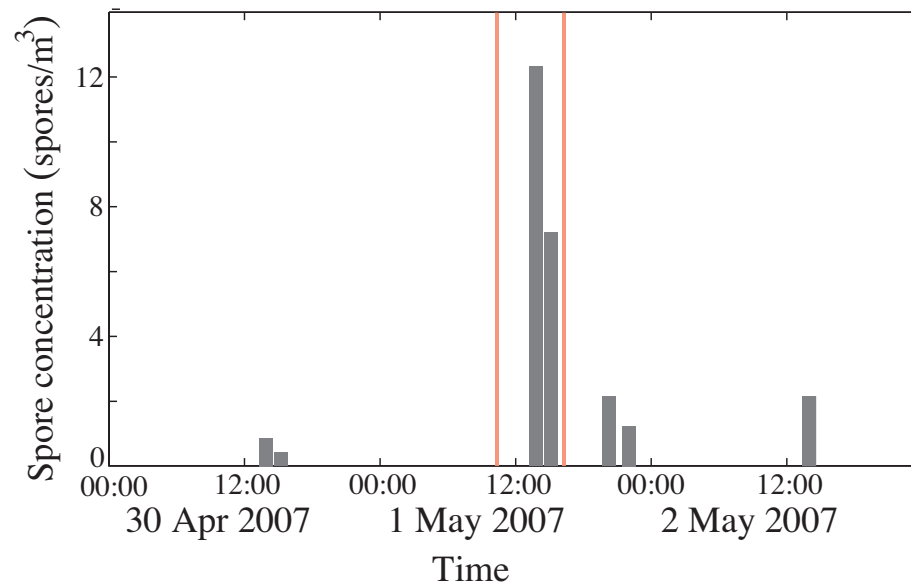
(a)



(b)



(c)



12:00 UTC 1 May 2007

15:00 UTC 1 May 2007

18:00 UTC 1 May 2007

Direct computation of coherent sets

- Take probabilistic point of view
- Partition phase space into **loosely coupled regions**

Coherent sets \approx “Leaky” regions with a long residence time³

phase space is divided into several invariant and almost-invariant sets.

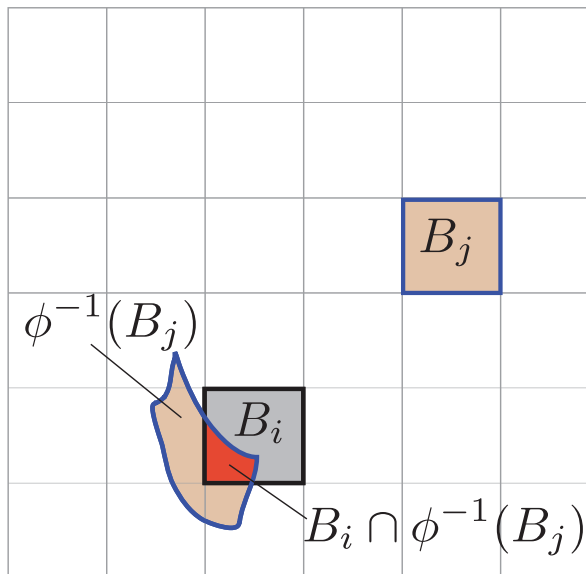
³work of Dellnitz, Junge, Deuffhard, Froyland, Schütte, et al

Direct computation of coherent sets

- Create box partition of phase space $\mathcal{B} = \{B_1, \dots, B_q\}$, with q large
- Consider a q -by- q **transition (Ulam) matrix**, P , for our dynamical system, where

$$P_{ij} = \frac{m(B_i \cap \phi^{-1}(B_j))}{m(B_i)},$$

the *transition probability* from B_i to B_j using $\phi = \phi_t^{t+T}$



- P approximates the flow map ϕ_t^{t+T} via a finite state Markov chain.

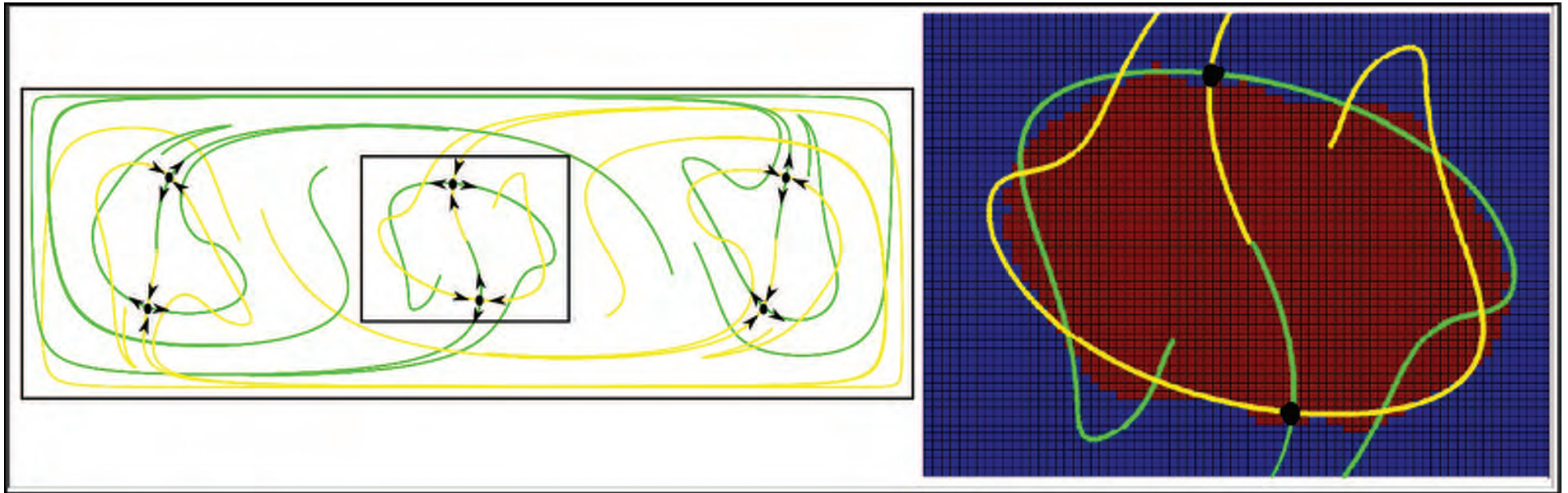
Direct computation of coherent sets

- A set B is called *almost-invariant* over the interval $[t, t + T]$ if

$$\rho(B) = \frac{m(B \cap \phi^{-1}(B))}{m(B)} \approx 1.$$

- Can maximize value of ρ over all possible combinations of sets $B \in \mathcal{B}$.
- In practice, AISs or relatedly, almost-cyclic sets (ACSs), identified via **eigenvectors** (of eigenvalues with $|\lambda| \approx 1$) of P or graph-partitioning

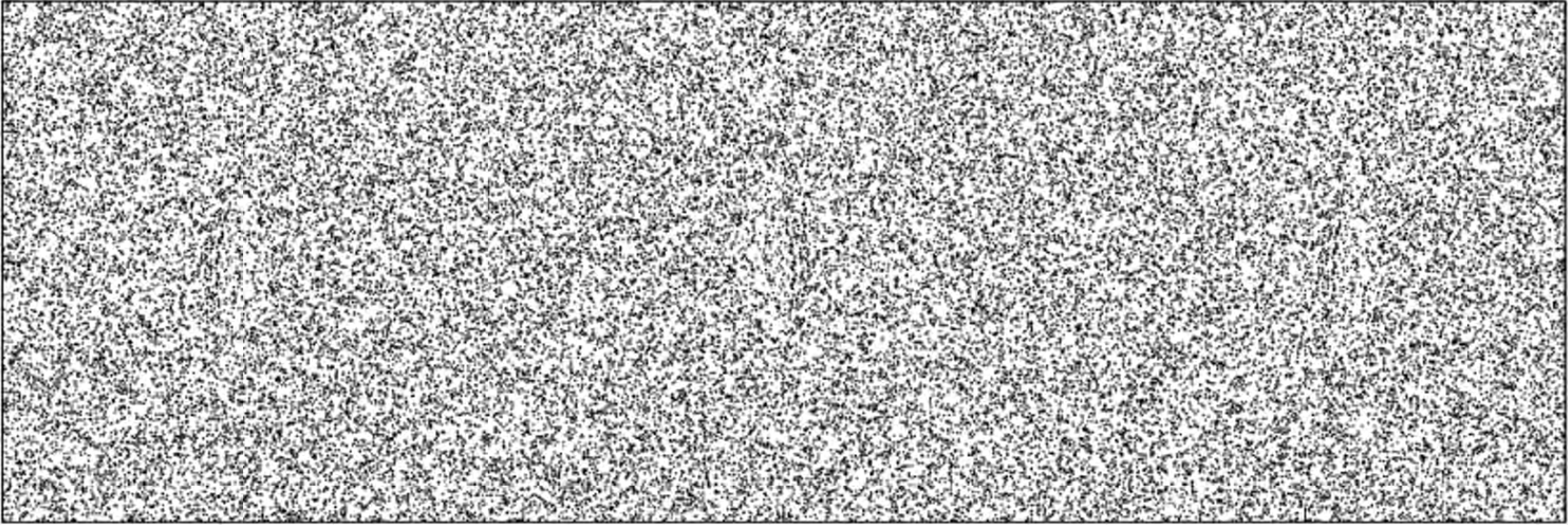
Identifying coherent sets



- Consider lid-driven cavity flow system with system parameter τ_f
- For $\tau_f > 1$, periodic points and manifolds exist
- Coherent set boundaries are manifolds of periodic points
- Known previously⁴ and applies to more general objects than periodic points, i.e. normally hyperbolic invariant manifolds (NHIMs)

⁴Dellnitz, Junge, Lo, Marsden, Padberg, Preis, Ross, Thiere [2005] Phys. Rev. Lett.; Dellnitz, Junge, Koon, Lekien, Lo, Marsden, Padberg, Preis, Ross, Thiere [2005] Int. J. Bif. Chaos

Identifying coherent sets

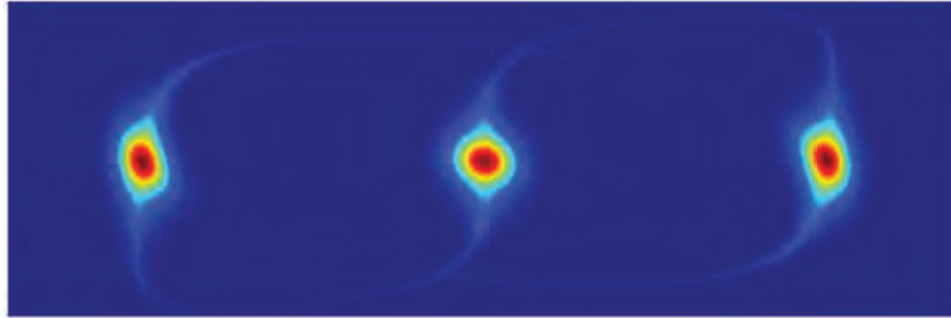


Poincaré section for $\tau_f < 1 \Rightarrow$ no obvious structure!

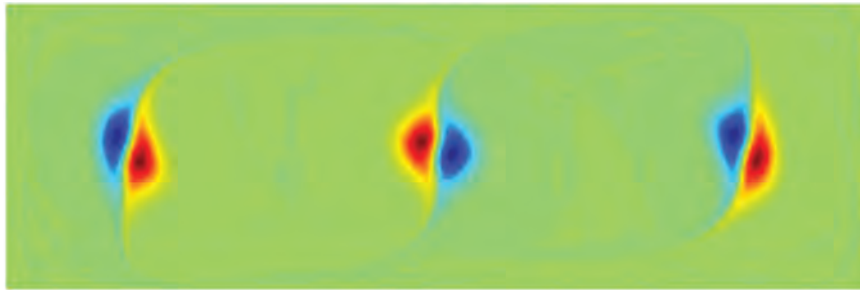
- For $\tau_f < 1$, no periodic orbits of low period known
- Is the phase space featureless?
- Consider transition matrix $P_t^{t+\tau_f}$ induced by Poincaré map $\phi_t^{t+\tau_f}$

Identifying coherent sets

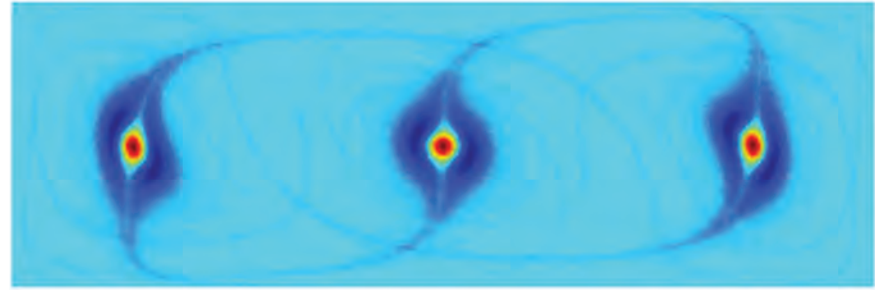
Top eigenvectors for $\tau_f = 0.99$ reveal hierarchy of phase space structures



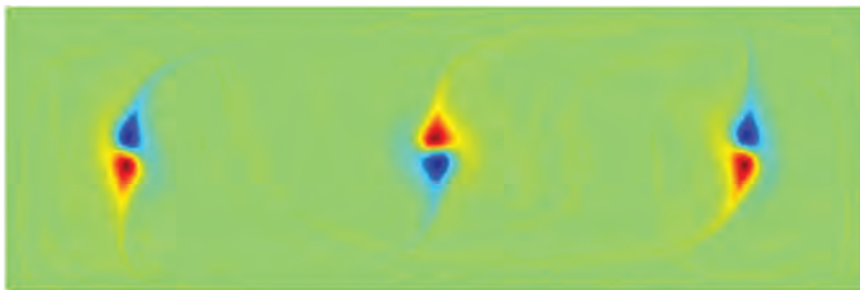
ν_2



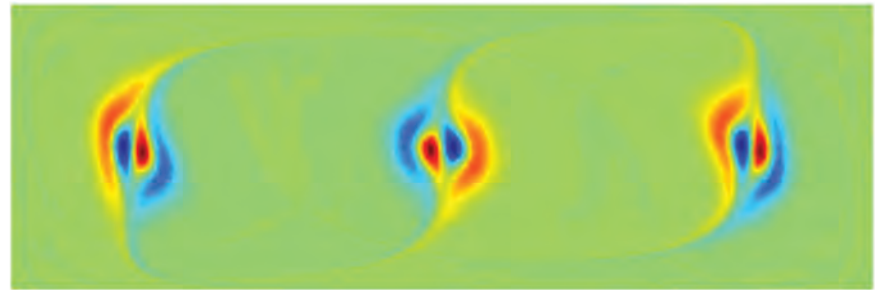
ν_3



ν_4

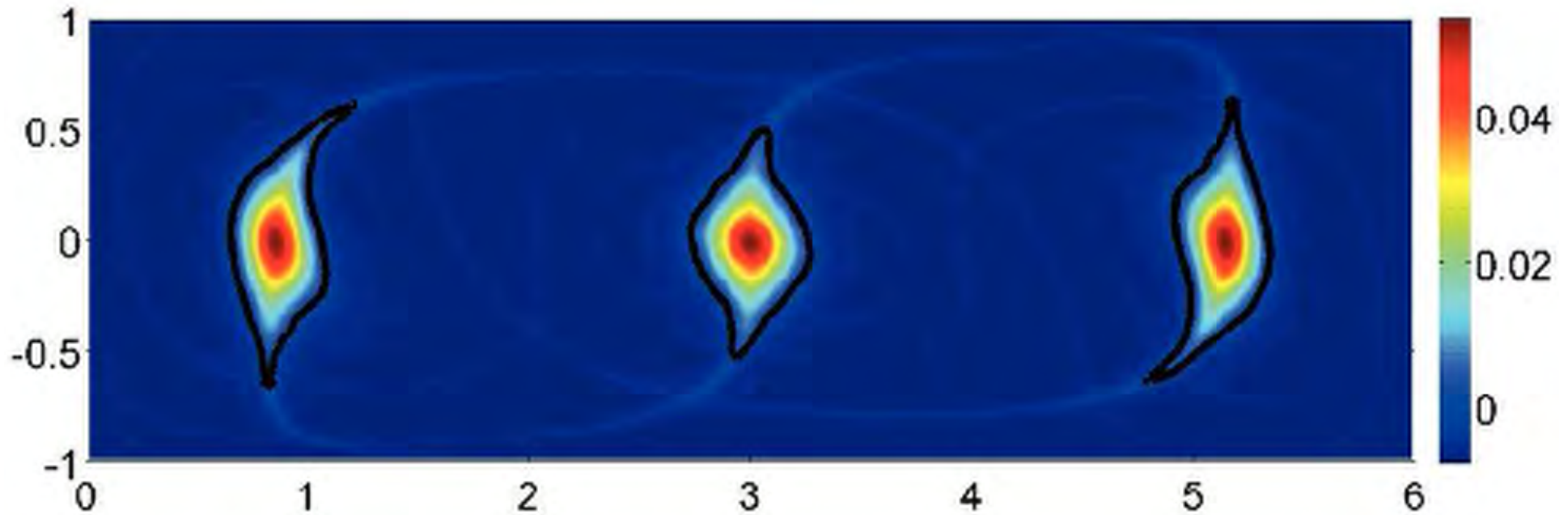


ν_5



ν_6

Identifying coherent sets



The zero contour (black) is the boundary between the two almost-invariant sets.

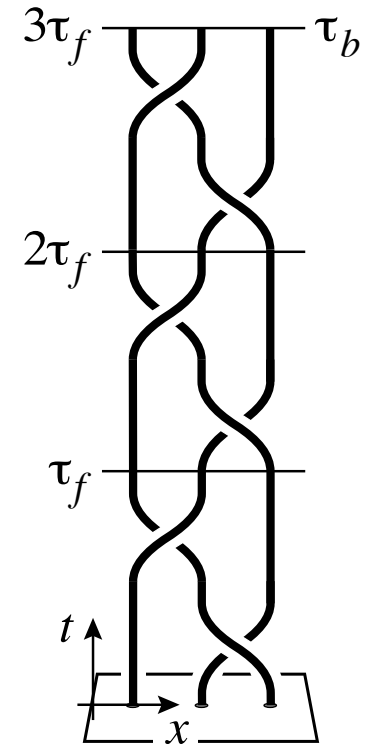
- Three almost-cyclic coherent sets (ACSs) of period 3
- ACSs effectively replace compact region bounded by saddle manifolds
- Also: we see a **dynamical remnant of the global ‘stable and unstable manifolds’ of the saddle points**, even there are no more saddle points

Identifying coherent sets

Almost-cyclic sets stirring the surrounding fluid like ‘ghost rods’
— **works even when periodic orbits are absent!**

Movie shown is second eigenvector for $P_t^{t+\tau_f}$ for $t \in [0, \tau_f)$

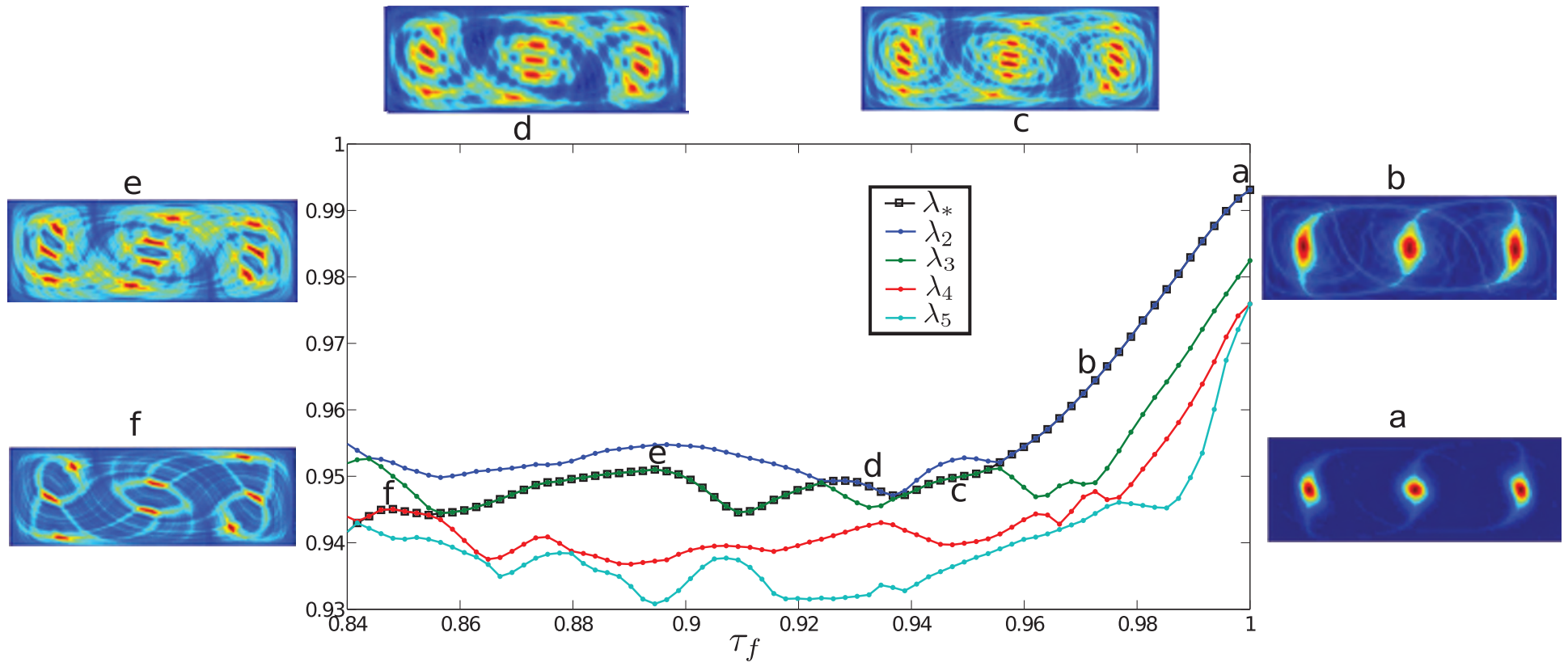
Identifying coherent sets



Braid of ACSs gives lower bound of entropy via Thurston-Nielsen

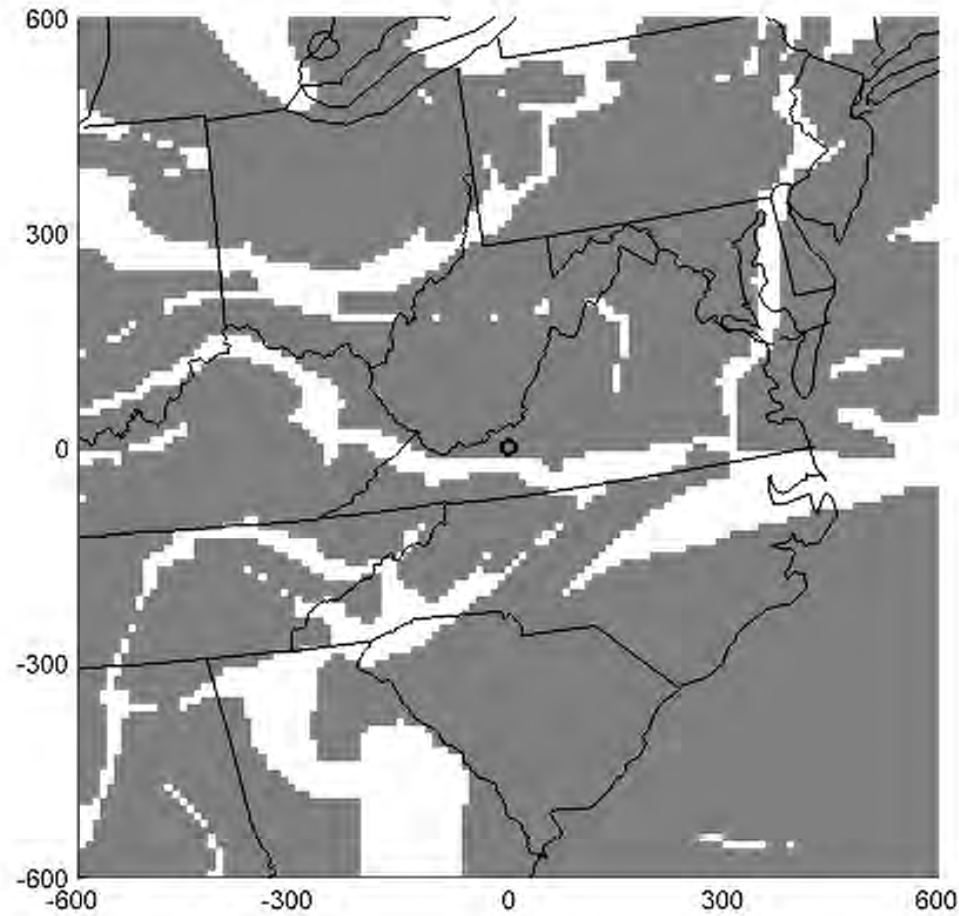
- One only needs approximately cyclic blobs of fluid
- Even though the theorems require exactly periodic points!
- Stremler, Ross, Grover, Kumar [2011] Phys. Rev. Lett.

Eigenvalues/eigenvectors vs. bifurcation parameter



Movie shows change in eigenvector along branch marked with '-□-' above (from a to f), as τ_f decreases \Rightarrow

Coherent sets in the atmosphere



- Coherent sets during 24 hours starting 09:00 1 May 2007

Final words on coherent sets & transport

- What are the robust descriptions of coherent sets and transport which work in aperiodic, finite-time settings?
- Methods for finding boundaries of coherent sets and coherent sets themselves are fruitful for illuminating structure and transport.
- Lobe dynamics, finite-time symbolic dynamics...
- In analogy with point vortices, can we find equations of motion for generalized coherent sets and their influence on each other?

The End

For papers, movies, etc., visit:

www.shaneross.com

Main Papers:

- Ross & Tallapragada [2011] Detecting and exploiting chaotic transport in mechanical systems, *Applications of Chaos and Non-linear Dynamics in Engineering - Vol. 2*, Springer.
- Tallapragada & Ross [2011] A geometric and probabilistic description of coherent sets, *under review*.
- Tallapragada, Ross, Schmale [2011] Lagrangian coherent structures are associated with fluctuations in airborne microbial populations. *Chaos* 21, 033122.
- Stremmler, Ross, Grover, Kumar [2011] Topological chaos and periodic braiding of almost-cyclic sets. *Physical Review Letters* 106, 114101.
- Senatore & Ross [2011] Detection and characterization of transport barriers in complex flows via ridge extraction of the finite time Lyapunov exponent field, *International Journal for Numerical Methods in Engineering* 86, 1163.
- Lekien & Ross [2010] The computation of finite-time Lyapunov exponents on unstructured meshes and for non-Euclidean manifolds. *Chaos* 20, 017505.